

Letter

Observation of a Be double-Lambda hypernucleus in the J-PARC E07 experiment

H. Ekawa^{1,2,*}, K. Agari³, J. K. Ahn⁴, T. Akaishi⁵, Y. Akazawa³, S. Ashikaga^{1,2}, B. Bassalleck⁶, S. Bleser⁷, Y. Endo⁸, Y. Fujikawa¹, N. Fujioka⁹, M. Fujita⁹, R. Goto⁸, Y. Han¹⁰, S. Hasegawa², T. Hashimoto², S. H. Hayakawa^{2,5}, T. Hayakawa⁵, E. Hayata¹, K. Hicks¹¹, E. Hirose³, M. Hirose¹, R. Honda⁹, K. Hoshino⁸, S. Hoshino⁵, K. Hosomi², S. H. Hwang¹², Y. Ichikawa², M. Ichikawa^{1,13}, M. Ieiri³, K. Imai², K. Inaba¹, Y. Ishikawa⁹, A. Iskenderi⁵, H. Ito⁸, K. Ito¹⁴, W. S. Jung⁴, S. Kanatsuki¹, H. Kanauchi⁹, A. Kasagi⁸, T. Kawai¹⁵, M. H. Kim⁴, S. H. Kim⁴, S. Kinbara^{2,8}, R. Kiuchi¹⁶, H. Kobayashi⁸, K. Kobayashi⁵, T. Koike⁹, A. Koshikawa¹, J. Y. Lee¹⁷, J. W. Lee⁴, T. L. Ma¹⁸, S. Y. Matsumoto^{1,13}, M. Minakawa³, K. Miwa⁹, A. T. Moe¹⁹, T. J. Moon¹⁷, M. Moritsu³, Y. Nagase⁸, Y. Nakada⁵, M. Nakagawa⁵, D. Nakashima⁸, K. Nakazawa⁸, T. Nanamura^{1,2}, M. Naruki^{1,2}, A. N. L. Nyaw⁸, Y. Ogura⁹, M. Ohashi⁸, K. Oue⁵, S. Ozawa⁹, J. Pochodzalla^{7,20}, S. Y. Ryu²¹, H. Sako², Y. Sasaki⁹, S. Sato², Y. Sato³, F. Schupp⁷, K. Shirotori²¹, M. M. Soe²², M. K. Soe⁸, J. Y. Sohn²³, H. Sugimura²⁴, K. N. Suzuki^{1,2}, H. Takahashi³, T. Takahashi³, Y. Takahashi¹, T. Takeda¹, H. Tamura^{2,9}, K. Tanida², A. M. M. Theint⁸, K. T. Tint⁸, Y. Toyama⁹, M. Ukai³, E. Umezaki¹, T. Watabe¹⁴, K. Watanabe¹, T. O. Yamamoto², S. B. Yang⁴, C. S. Yoon²³, J. Yoshida², M. Yoshimoto⁸, D. H. Zhang¹⁸, and Z. Zhang¹⁸

¹Department of Physics, Kyoto University, Kyoto 606-8502, Japan

²Advanced Science Research Center, Japan Atomic Energy Agency, Tokai 319-1195, Japan

³Institute of Particle and Nuclear Study (IPNS), High Energy Accelerator Research Organization (KEK), Tsukuba 305-0801, Japan

⁴Department of Physics, Korea University, Seoul 02841, Korea

⁵Department of Physics, Osaka University, Toyonaka 560-0043, Japan

⁶Department of Physics and Astronomy, University of New Mexico, Albuquerque, NM 87131, USA

⁷Helmholtz Institute Mainz, 55099 Mainz, Germany

⁸Physics Department, Gifu University, Gifu 501-1193, Japan

⁹Department of Physics, Tohoku University, Sendai 980-8578, Japan

¹⁰Institute of Nuclear Energy Safety Technology, Chinese Academy of Sciences, Hefei 230031, China

¹¹Department of Physics & Astronomy, Ohio University, Athens, OH 45701, USA

¹²Korea Research Institute of Standards and Science, Daejeon 34113, Korea

¹³RIKEN Cluster for Pioneering Research, Wako 351-0198, Japan

¹⁴Department of Physics, Nagoya University, Nagoya 464-8601, Japan

¹⁵RIKEN Nishina Center, Wako 351-0198, Japan

¹⁶Institute of High Energy Physics, Beijing 100049, China

¹⁷Department of Physics, Seoul National University, Seoul 08826, Korea

¹⁸Institute of Modern Physics, Shanxi Normal University, Linfen 041004, China

¹⁹Department of Physics, Lashio University, Buda Lane, Lashio 06301, Myanmar

²⁰Institut für Kernphysik, Johannes Gutenberg-Universität, 55099 Mainz, Germany

²¹Research Center for Nuclear Physics, Osaka University, Osaka 567-0047, Japan

²²Department of Physics, University of Yangon, Yangon 11041, Myanmar

²³Research Institute of Natural Science, Gyeongsang National University, Jinju 52828, Korea

²⁴Accelerator Laboratory, High Energy Accelerator Research Organization (KEK), Tsukuba 305-0801, Japan

*E-mail: ekawa@scphys.kyoto-u.ac.jp

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A double- Λ hypernucleus, ${}_{\Lambda\Lambda}\text{Be}$, was observed by the J-PARC E07 Collaboration in nuclear emulsions tagged by the (K^-, K^+) reaction. This event was interpreted as the production and decay of ${}_{\Lambda\Lambda}^{10}\text{Be}$, ${}_{\Lambda\Lambda}^{11}\text{Be}$, or ${}_{\Lambda\Lambda}^{12}\text{Be}^*$ via Ξ^- capture in ${}^{16}\text{O}$. By assuming capture in the atomic 3D state, the binding energies of two Λ hyperons ($B_{\Lambda\Lambda}$) of these double- Λ hypernuclei are obtained to be 15.05 ± 0.11 MeV, 19.07 ± 0.11 MeV, and 13.68 ± 0.11 MeV, respectively. Based on the kinematic fitting, ${}_{\Lambda\Lambda}^{11}\text{Be}$ is the most likely explanation for the observed event.
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1. Introduction A complete understanding of the baryon–baryon interaction requires (at least) the consideration of all octet baryons within the $\text{SU}(3)_f$ group. Compared to baryon pairs with zero or only one strange baryon, experimental information in the $S = -2$ sector is still very scarce. In particular, it is difficult to study the interaction between two hyperons by scattering experiments due to their short lifetimes. Therefore, double- Λ hypernuclei, which include two Λ hyperons in a nucleus, have been investigated. The $\Lambda\Lambda$ interaction is expressed in terms of $\Delta B_{\Lambda\Lambda}$, which can be deduced from the mass of a double- Λ hypernucleus and which is defined as

$$\Delta B_{\Lambda\Lambda}({}_{\Lambda\Lambda}^AZ) = B_{\Lambda\Lambda}({}_{\Lambda\Lambda}^AZ) - 2B_{\Lambda}({}_{\Lambda}^{A-1}Z). \quad (1)$$

Here B_{Λ} and $B_{\Lambda\Lambda}$ represent the binding energies of a Λ hyperon in single- Λ hypernuclei and two Λ hyperons in double- Λ hypernuclei, respectively. Emulsion detectors allow the detection of the weak decay products from double- Λ hypernuclei with sub- μm resolution, thus providing the most precise reconstruction of double- Λ hypernucleus masses as of today.

In the past, several experiments have successfully discovered double- Λ hypernuclear decays in nuclear emulsions [1–4]. The most impressive results were collected by the KEK-PS E373 experiment. Among seven double- Λ hypernuclear events, an event called “NAGARA” was uniquely identified as ${}_{\Lambda\Lambda}^6\text{He}$ [3]. From this event, the $\Lambda\Lambda$ interaction, and especially its s -wave (1S_0) interaction, was found to be weakly attractive ($\Delta B_{\Lambda\Lambda} = 0.67 \pm 0.17$ MeV) [4]. In order to understand $\Lambda\Lambda$ and the related interaction systematically, more double- Λ hypernuclei need to be uniquely identified.

J-PARC E07 is an upgraded counter-emulsion hybrid experiment aiming to identify about a factor of 10 more double- Λ hypernuclei as compared to the E373 experiment [5]. It is expected to detect approximately 100 double- Λ hypernuclear events among 1×10^4 Ξ^- stopping events. This substantially improved statistics will allow us to explore not only the $\Lambda\Lambda$ s -wave interaction but also ΞN – $\Lambda\Lambda$ mixing, the structure of the core nucleus, etc.

Beam exposure of E07 was carried out in 2016 and 2017. A total of 118 modules produced from 2.1 tons of emulsion gel were exposed to 1.13×10^{11} particles of the K^- beam. An impressive double- Λ hypernuclear event, “MINO”, was observed after scanning 30% of all modules. In this letter, an interpretation of this event is discussed.

2. Experimental setup Double- Λ hypernuclei were searched for by detecting Ξ^- stopping events in an emulsion module. A Ξ^- hyperon generated in the quasi-free (K^-, K^+) reaction in

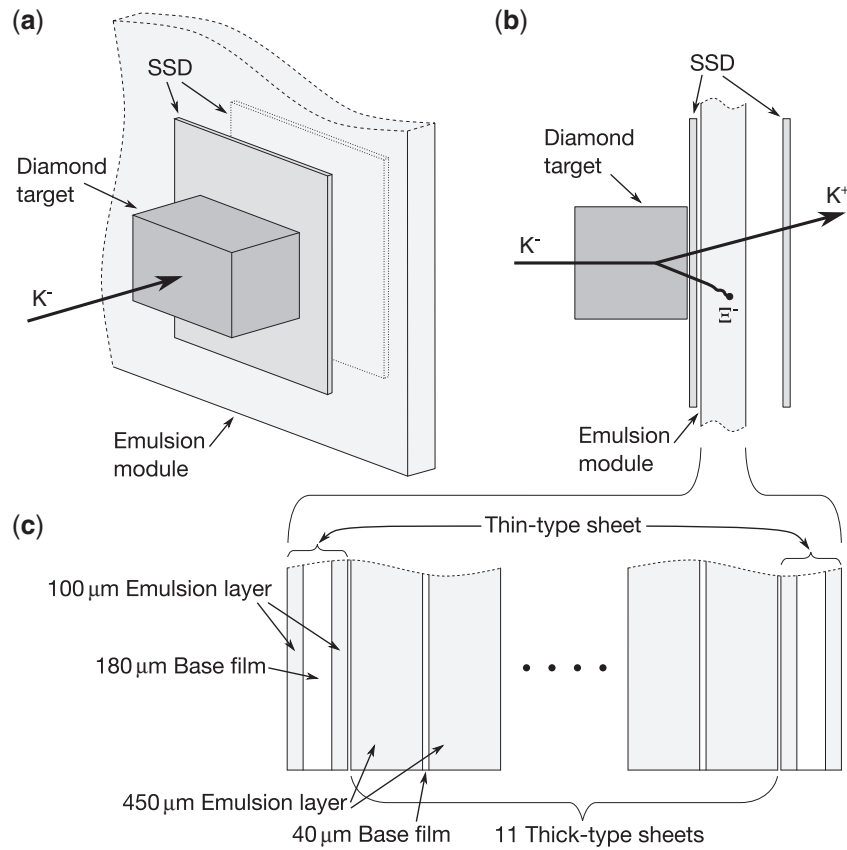


Fig. 1. Schematic view of the experimental setup around a diamond target. The target size was 50 mm (W) $\times$ 30 mm (H) $\times$ 30 mm (T).

a diamond target with 30 mm thickness (9.83 g/cm^2) was injected into the module and subsequently slowed down and captured in the atomic orbit of a nucleus in the emulsion material. Double- Λ hypernuclei are produced by the interaction between the Ξ^- hyperon and nucleus with a probability of a few %. In order to trace Ξ^- tracks to the module and detect particle tracks from the module, two silicon strip detector (SSD) sets were installed so as to sandwich the module. The experimental setup around the target is shown in Fig. 1.

The J-PARC E07 experiment was performed at the K1.8 beam line in the J-PARC Hadron Experimental Facility with K^- beams of $1.8 \text{ GeV}/c$ momentum. This momentum was chosen to maximize the Ξ^- stopping yield in the emulsion. The typical intensity and purity of the K^- beam were 2.8×10^5 particles per spill of 2.0 s duration every 5.52 s and 82%, respectively. The momenta of the incoming K^- and outgoing K^+ mesons were analyzed by corresponding magnetic spectrometers, the beam-line spectrometer [6] and the KURAMA spectrometer, respectively. The momentum resolution of each spectrometer was $\Delta p/p = 3.3 \times 10^{-4}$ (FWHM) and $\Delta p/p = 2.7 \times 10^{-2}$ (FWHM), respectively. The acceptance of the KURAMA spectrometer was 280 msr.

Each emulsion module consisted of 11 thick-type sheets sandwiched between two thin-type sheets with an area of 345 mm (W) $\times$ 350 mm (H) (Fig. 1(c)). The thin-type sheets had emulsion layers with a thickness of $100 \mu\text{m}$ on both sides of $180 \mu\text{m}$ polystyrene base film and were used to connect tracks to the SSDs, because they have high deformation tolerance and thus good angular resolution. The thick-type sheets had $450 \mu\text{m}$ thick layers on both sides of $40 \mu\text{m}$ polystyrene base film. The

emulsion layers were made of “Fuji GIF” emulsion gel produced by the FUJIFILM Corporation. The 13 sheets were packed all together in a stainless case and fixed tightly by vacuum pumping.

The module was moved in the beam spill-off period to keep the beam particle density less than 1×10^4 particles/mm² in order to maintain good efficiency for automated image tracking. A typical exposure cycle was 5 hr for one module.

Each SSD had a four-layer configuration (XYXY) with a strip pitch of 50 μ m. The thickness of a silicon sensor was 0.3 mm. The resolutions of the SSDs were estimated to be 15 μ m (σ) for position and 20 mrad (σ) for angle.

3. Analysis Tracks of Ξ^- were kinematically identified with the SSD by tagging the (K^-, K^+) reaction. The candidate tracks of Ξ^- hyperons were constructed from hits in four layers checking the consistency with the “ p ”(K^-, K^+) Ξ^- kinematics. In order to select Ξ^- hyperons with high stopping probability, Ξ^- candidates with energy deposits more than six times larger than the minimum ionizing particles in the SSD were selected. The downstream SSD was used to reject Ξ^- candidates that penetrated the module without nuclear interactions. A relative position between the module and the SSDs was calibrated using $\bar{p}$ beam through events with a beam-pattern matching method. By calibrating the four corners of the module, it was corrected with 20 μ m accuracy. The prediction accuracy of the Ξ^- incident position was estimated to be about 50 μ m for the thin-type sheet.

From the sets of the predicted positions and angles of the Ξ^- hyperons based on the SSD hits, the tracks were traced through the emulsion with automated microscope systems [7]. First, the most upstream sheet of the module was scanned to find Ξ^- tracks in an area of about 200 μ m $\times$ 200 μ m for each Ξ^- hit position predicted in the previous step of analysis. Then, the Ξ^- tracks found in the first thin-type sheet were traced downstream through several emulsion sheets. When the tracing by the microscope system detected the end point of the track, the system took photographs around the stopping point and these were checked by human eyes.

For kinematic analysis, range–energy calibration and shrinkage correction were necessary and performed for each emulsion sheet by α tracks with a monochromatic energy of 8.784 MeV from the decay of ^{212}Po existing in the emulsion. The α track can be identified in the thorium series isotopes because it has the largest kinetic energy, i.e., the longest range. Such α decay chains were searched for around the observed event by using the so-called overall scanning method [8]. One hundred α tracks were scanned for calibration. Since the emulsion layers were shrunk along the beam direction due to photographic development, the ranges of particles needed to be corrected for the shrinkage effect. The mean range of α tracks and the shrinkage factor were obtained to be 50.77 ± 0.12 μ m and 1.98 ± 0.02 , respectively. The relation between the ranges and kinetic energies of charged particles was obtained by the range–energy formula given by Barkas et al. [9,10]. The density of the emulsion sheet was determined to be 3.486 ± 0.013 g/cm³. The ranges of particles in polystyrene base films and SSDs were converted into corresponding emulsion ranges considering the energy loss ratio.

4. Interpretation of the MINO event A new double- Λ hypernuclear event was observed in the seventh sheet of a module. An overlaid photograph and a schematic drawing of the event are shown in Fig. 2. We called this event “MINO”¹.

The Ξ^- hyperon came to rest at vertex A, from which three charged particles (#1, #3, and #4) were emitted. The particle of track #1 decayed to three charged particles (#2, #5, and #6) at vertex B. The particle of track #2 decayed again to three charged particles (#7, #8, and #9) at vertex C. The

¹ The name of the southern part of Gifu prefecture, Japan, where the event was found.

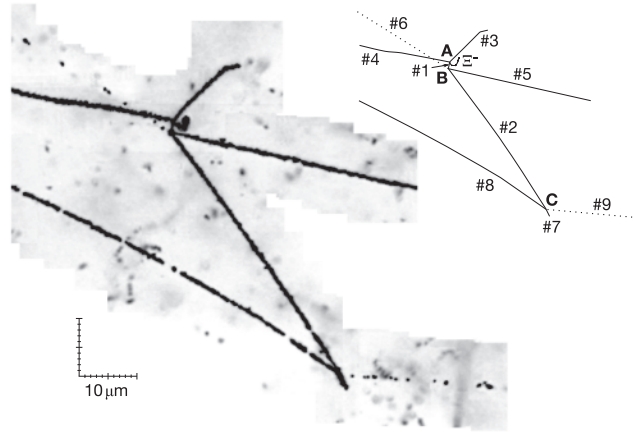


Fig. 2. A photograph of the MINO event and its schematic drawing. The overlaid photograph is made by patching focused regions. Tracks #4, #5, #6, #8, and #9 are not fully shown in this photograph because these tracks are too long to be presented.

Table 1. Ranges and angles of the tracks of the MINO event. The zenith and azimuthal angles are presented in columns θ and ϕ .

Vertex	Track ID	Range [μm]	θ [degree]	ϕ [degree]	Comment
A	#1	2.1 ± 0.2	83.7 ± 8.9	256.1 ± 5.3	double- Λ hypernucleus
	#3	17.5 ± 0.2	121.9 ± 1.9	48.2 ± 1.3	
	#4	65.7 ± 0.5	41.7 ± 1.7	166.7 ± 2.1	
B	#2	50.6 ± 0.3	90.2 ± 2.0	306.3 ± 1.3	single- Λ hypernucleus
	#5	122.1 ± 0.2	61.4 ± 1.8	347.0 ± 1.5	
	#6	$> 23\ 170$	106.2 ± 0.6	147.7 ± 0.4	
C	#7	5.0 ± 0.2	31.1 ± 2.8	297.0 ± 4.0	stopped in the SSD
	#8	116.7 ± 0.2	100.3 ± 1.9	144.2 ± 1.3	
	#9	> 7378	147.4 ± 0.3	355.7 ± 0.5	

measured ranges and angles are summarized in Table 1. If the Ξ^- hyperon was captured in a heavy nucleus such as Ag or Br, a short track like #3 with a range of less than $32\ \mu\text{m}$ could not be emitted due to the Coulomb barrier [2]. Therefore, we have concluded that the Ξ^- hyperon was captured in a light nucleus such as ^{12}C , ^{14}N , or ^{16}O . The particles of tracks #6 and #9 escaped from the module into the downstream SSD after passing several emulsion sheets. These tracks could be connected to the SSD by extrapolating the tracks at the exit point from the last emulsion sheet. The particle of track #6 was found to be stopped in the fourth layer of the SSD and #9 penetrated all SSD layers. The ranges of #6 and #9 in the SSD were $4500 \pm 200\ \mu\text{m}$ and $2200 \pm 20\ \mu\text{m}$ in emulsion equivalent, respectively.

We began by checking vertex C. The coplanarity is defined as $(\vec{r}_1 \times \vec{r}_2) \cdot \vec{r}_3$, where $\vec{r}_i$ is a unit vector of a track angle. Three charged particles were emitted with a coplanarity of 0.001 ± 0.043 . This shows that three particles were emitted in a plane; thus, neutron emission is unlikely. The possibility of neutron emission is discussed at the end of this section. From all nuclide combinations for both mesonic and non-mesonic decays of known single- Λ hypernuclei, possible decay modes were selected using the following criteria. (1) An angular difference between #9 and the momentum sum of #7 and #8 should be back-to-back with 3σ confidence. (2) Momenta and energies should be

Table 2. Possible decay modes at vertex C in the case of no neutron emission. Candidates that are accepted by the angular constraint and the conservation of momentum and energy in the 3σ cut condition are listed. The χ^2 value and the total range of #9 were obtained from the kinematic fitting [11].

Single- Λ hypernucleus (#2)		#7	#8	#9	χ^2	Range (#9) [μm]	Comment
${}^4_{\Lambda}\text{He}$	$\rightarrow$	${}^3\text{He}$	p	π^-	33.1	16 800	rejected
${}^5_{\Lambda}\text{He}$	$\rightarrow$	${}^4\text{He}$	p	π^-	5.23	16 270	
${}^8_{\Lambda}\text{Li}$	$\rightarrow$	${}^6\text{Li}$	d	π^-	93.6	7906	rejected
${}^9_{\Lambda}\text{Li}$	$\rightarrow$	${}^7\text{Li}$	d	π^-	105	10 660	rejected

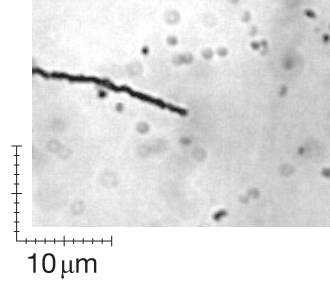


Fig. 3. A photograph of the end point of track #5.

conserved with 3σ by applying the kinematic fitting with a degree of freedom (DOF) of 3 [11]. Here, the range of #9 was parameterized to conserve the total momentum and reconstruct the mass of a single- Λ hypernucleus. Possible decay modes at vertex C are listed in Table 2. When the chi-squared value of the kinematic fitting was larger than 14.2, such decay modes were rejected. In this case, the p-value for the fitting is evaluated as 0.27%, which corresponds to the 3σ cut condition. Taking this into account, a possible candidate for #2 was identified as ${}^5_{\Lambda}\text{He}$ in the case of no neutron emission. The lower limit of the range of #9 was obtained as 7378 (in emulsion) + 2200 (in SSD) μm considering the track length in the SSD. The interpretation of ${}^5_{\Lambda}\text{He}$ is consistent with this requirement.

Next, we checked vertex B. The particle of track #1 decayed to three charged particles including a very thin track (#6). The range of #6 was measured as 23 170 (in the emulsion) + 4500 (in the SSD) μm . If the particle of track #6 is π^- , the total visible energy by decay daughters (#2, #5, and #6) is at least 47.7 MeV. Since this energy is larger than the Q value of any π^- mesonic decay mode, the possibility of #6 being π^- was rejected. Thus, the charge of #1 should be more than three. The maximum charge of #1 is five by assuming that Ξ^- was captured in ${}^{16}\text{O}$ because three charged particles were emitted at vertex A. Therefore, the nuclide of #1 is ${}_{\Lambda\Lambda}\text{Be}$ or ${}_{\Lambda\Lambda}\text{B}$.

The coplanarity of vertex B was calculated to be 0.007 ± 0.019 . Among the decay modes with ${}^5_{\Lambda}\text{He}$ as #2 without neutron emission, only the following decay satisfied kinematical consistency:

$${}^{13}_{\Lambda\Lambda}\text{B} \rightarrow {}^5_{\Lambda}\text{He} + {}^6\text{He} + d.$$

However, this decay mode was also rejected because there was no electron track associated with the end point of track #5 as seen in Fig. 3, even though ${}^6\text{He}$ should decay to ${}^6\text{Li} + e^- + \bar{\nu}$ with a half-life of 806.7 ms [12].

Thus, neutron(s) should be emitted at vertex B although the coplanarity is so small. Regarding decay modes with neutron(s) at vertex B, all nuclide combinations were checked for charged particles. In the kinematic analysis, the range of #6 was calculated by assuming a double- Λ hypernucleus

Table 3. Possible decay modes at vertex B.

Double- Λ hypernucleus (#1)		#2	#5	#6
${}^9_{\Lambda\Lambda}\text{Be}$	$\rightarrow$	${}^5_{\Lambda}\text{He}$	p	$p \quad 2n$
${}^{10}_{\Lambda\Lambda}\text{Be}$	$\rightarrow$	${}^5_{\Lambda}\text{He}$	(p, d)	$p \quad (3n, 2n)$
${}^{11}_{\Lambda\Lambda}\text{Be}$	$\rightarrow$	${}^5_{\Lambda}\text{He}$	(p, d, t)	$p \quad (4n, 3n, 2n)$
${}^{12}_{\Lambda\Lambda}\text{Be}$	$\rightarrow$	${}^5_{\Lambda}\text{He}$	(p, d, t)	$p \quad (5n, 4n, 3n)$
${}^{13}_{\Lambda\Lambda}\text{Be}$	$\rightarrow$	${}^5_{\Lambda}\text{He}$	(p, d, t)	$p \quad (6n, 5n, 4n)$
${}^{11}_{\Lambda\Lambda}\text{B}$	$\rightarrow$	${}^5_{\Lambda}\text{He}$	${}^3\text{He}$	$p \quad 2n$
${}^{12}_{\Lambda\Lambda}\text{B}$	$\rightarrow$	${}^5_{\Lambda}\text{He}$	$({}^3\text{He}, {}^4\text{He})$	$p \quad (3n, 2n)$
${}^{12}_{\Lambda\Lambda}\text{B}$	$\rightarrow$	${}^5_{\Lambda}\text{He}$	${}^4\text{He}$	$d \quad n$
${}^{13}_{\Lambda\Lambda}\text{B}$	$\rightarrow$	${}^5_{\Lambda}\text{He}$	$({}^3\text{He}, {}^4\text{He})$	$p \quad (4n, 3n)$
${}^{14}_{\Lambda\Lambda}\text{B}$	$\rightarrow$	${}^5_{\Lambda}\text{He}$	${}^4\text{He}$	$p \quad 4n$
${}^{15}_{\Lambda\Lambda}\text{B}$	$\rightarrow$	${}^5_{\Lambda}\text{He}$	${}^4\text{He}$	$p \quad 5n$

with $\Delta B_{\Lambda\Lambda} = 0$, where the missing momentum was carried by unobserved neutron(s). In the case of multiple neutron emissions, all neutrons were treated as having the same momentum. This setting gives the minimum kinetic energy of neutrons and the maximum kinetic energy of #6, which corresponds to the maximum range of #6. If the maximum range was not consistent with the measurement, those assignments were rejected. Since non-mesonic decays have large Q values, many decay modes remained for the case of #1 being ${}_{\Lambda\Lambda}\text{Be}$ or ${}_{\Lambda\Lambda}\text{B}$ nuclides, as summarized in Table 3.

Finally, we checked vertex A, where three tracks were observed. All nuclide combinations for #1 to be ${}_{\Lambda\Lambda}\text{Be}$ or ${}_{\Lambda\Lambda}\text{B}$ were checked. In case of the decay with neutron emission, the momentum of the neutron(s) was assumed to be the missing momentum. In the case of more than one neutron emission, only a lower limit of $\Delta B_{\Lambda\Lambda}$ could be obtained. Possible decay modes are listed in Table 4. However, neutron emission was unlikely because the coplanarity of vertex A was calculated to be 0.000 ± 0.099 . Additionally, $\Delta B_{\Lambda\Lambda} - B_{\Xi^-}$ should not be a large value considering the result of the NAGARA event.

From the above considerations, the candidates for the production and decay modes are as follows:

$$\begin{aligned}
 {}^{16}\text{O} + \Xi^- &\rightarrow ({}^{10}_{\Lambda\Lambda}\text{Be}, {}^{11}_{\Lambda\Lambda}\text{Be}, {}^{12}_{\Lambda\Lambda}\text{Be}) + {}^4\text{He} + (t, d, p), \\
 &\hookrightarrow {}^5_{\Lambda}\text{He} + (p, d, t) + p + xn, \\
 &\hookrightarrow {}^4\text{He} + p + \pi^-.
 \end{aligned}$$

The nuclide of the double hypernucleus was uniquely identified as a ${}_{\Lambda\Lambda}\text{Be}$. The $B_{\Lambda\Lambda}$ and $\Delta B_{\Lambda\Lambda}$ values depend on the Ξ^- binding energy (B_{Ξ^-}). If we assume that the Ξ^- hyperon was captured in the atomic 3D state of ${}^{16}\text{O}$ with the theoretically estimated B_{Ξ^-} value of 0.23 MeV [13], $B_{\Lambda\Lambda}$ ($\Delta B_{\Lambda\Lambda}$) for each decay mode are obtained to be 15.05 ± 0.11 MeV (1.63 ± 0.14 MeV), 19.07 ± 0.11 MeV (1.87 ± 0.37 MeV), and 13.68 ± 0.11 MeV (-2.7 ± 1.0 MeV), for ${}^{10}_{\Lambda\Lambda}\text{Be}$, ${}^{11}_{\Lambda\Lambda}\text{Be}$, and ${}^{12}_{\Lambda\Lambda}\text{Be}$, respectively. These values are summarized in Table 5 together with their statistical and systematic errors. The statistical error was caused by the kinematic fitting and the systematic error was caused by the mass of the Ξ^- hyperon and $B_{\Lambda}({}^{A-1}_{\Lambda}Z)$. The masses of the Ξ^- and Λ hyperons were taken as 1321.71 ± 0.07 MeV and 1115.683 ± 0.006 MeV, respectively [14]. In this analysis, we took 8.2 ± 0.5 MeV for $B_{\Lambda}({}^{11}_{\Lambda}\text{Be})$ by a linear extrapolation from B_{Λ} of ${}_{\Lambda}\text{Be}$ isotopes with values of 5.16 ± 0.08 MeV (${}^7_{\Lambda}\text{Be}$) [15], 6.84 ± 0.05 MeV (${}^8_{\Lambda}\text{Be}$) [15], 6.71 ± 0.04 MeV (${}^9_{\Lambda}\text{Be}$) [15], and $8.60 \pm 0.07 \pm 0.16$ MeV (${}^{10}_{\Lambda}\text{Be}$) [16] because ${}^{11}_{\Lambda}\text{Be}$ has not yet been observed.

Table 4. Possible decay modes at Vertex A. The error of $\Delta B_{\Lambda\Lambda}$ shows the result of the kinematic fitting derived from our measurement. Candidates that have $\Delta B_{\Lambda\Lambda} - B_{\Xi^-} < 20$ MeV are listed.

Ξ^- capture	#1	#3	#4	$\Delta B_{\Lambda\Lambda} - B_{\Xi^-}$ [MeV]	
$^{16}\text{O} + \Xi^- \rightarrow$	$^{10}_{\Lambda\Lambda}\text{Be}$	^4He	t	1.40 ± 0.09	
$^{16}\text{O} + \Xi^- \rightarrow$	$^{11}_{\Lambda\Lambda}\text{Be}$	^4He	d	1.64 ± 0.08	
$^{16}\text{O} + \Xi^- \rightarrow$	$^{12}_{\Lambda\Lambda}\text{Be}$	^4He	p	-2.95 ± 0.08	
$^{14}\text{N} + \Xi^- \rightarrow$	$^{10}_{\Lambda\Lambda}\text{Be}$	p	p	$3n$	$> 14.18 \pm 0.64$
$^{14}\text{N} + \Xi^- \rightarrow$	$^{10}_{\Lambda\Lambda}\text{Be}$	p	d	$2n$	$> 15.94 \pm 1.22$
$^{14}\text{N} + \Xi^- \rightarrow$	$^{10}_{\Lambda\Lambda}\text{Be}$	d	p	$2n$	$> 12.18 \pm 0.72$
$^{14}\text{N} + \Xi^- \rightarrow$	$^{10}_{\Lambda\Lambda}\text{Be}$	d	d	n	16.29 ± 1.92
$^{14}\text{N} + \Xi^- \rightarrow$	$^{10}_{\Lambda\Lambda}\text{Be}$	t	p	n	7.08 ± 1.05
$^{14}\text{N} + \Xi^- \rightarrow$	$^{11}_{\Lambda\Lambda}\text{Be}$	p	p	$2n$	$> 10.60 \pm 0.93$
$^{14}\text{N} + \Xi^- \rightarrow$	$^{11}_{\Lambda\Lambda}\text{Be}$	p	d	n	17.43 ± 2.41
$^{14}\text{N} + \Xi^- \rightarrow$	$^{11}_{\Lambda\Lambda}\text{Be}$	d	p	n	10.30 ± 1.38
$^{14}\text{N} + \Xi^- \rightarrow$	$^{12}_{\Lambda\Lambda}\text{Be}$	p	p	n	10.21 ± 1.79
$^{16}\text{O} + \Xi^- \rightarrow$	$^{10}_{\Lambda\Lambda}\text{Be}$	^4He	p	$2n$	$> 9.35 \pm 0.47$
$^{16}\text{O} + \Xi^- \rightarrow$	$^{10}_{\Lambda\Lambda}\text{Be}$	^4He	d	n	7.73 ± 0.40
$^{16}\text{O} + \Xi^- \rightarrow$	$^{11}_{\Lambda\Lambda}\text{Be}$	^4He	p	n	4.42 ± 0.91
$^{16}\text{O} + \Xi^- \rightarrow$	$^{12}_{\Lambda\Lambda}\text{Be}$	^3He	p	n	19.30 ± 0.82
$^{16}\text{O} + \Xi^- \rightarrow$	$^{13}_{\Lambda\Lambda}\text{B}$	p	p	$2n$	$> 14.20 \pm 1.34$
$^{16}\text{O} + \Xi^- \rightarrow$	$^{13}_{\Lambda\Lambda}\text{B}$	d	p	n	16.55 ± 2.09
$^{16}\text{O} + \Xi^- \rightarrow$	$^{14}_{\Lambda\Lambda}\text{B}$	p	p	n	18.44 ± 2.65

Table 5. Result of $B_{\Lambda\Lambda}$ and $\Delta B_{\Lambda\Lambda}$ for the MINO event. The Ξ^- hyperon was assumed to be captured in the atomic 3D state of ^{16}O ($B_{\Xi^-} = 0.23$ MeV).

Nuclide	$B_{\Lambda\Lambda}$ [MeV]	$\Delta B_{\Lambda\Lambda}$ [MeV]	$B_{\Lambda}({}^{A-1}_{\Lambda}Z)$ [MeV]
$^{10}_{\Lambda\Lambda}\text{Be}$	$15.05 \pm 0.09 \pm 0.07$	$1.63 \pm 0.09 \pm 0.11$	6.71 ± 0.04 [15]
$^{11}_{\Lambda\Lambda}\text{Be}$	$19.07 \pm 0.08 \pm 0.07$	$1.87 \pm 0.08 \pm 0.36$	$8.60 \pm 0.07 \pm 0.16$ [16]
$^{12}_{\Lambda\Lambda}\text{Be}$	$13.68 \pm 0.08 \pm 0.07$	$(-2.7 \pm 0.08 \pm 1.0)$	8.2 ± 0.5 (extrapolation)

Table 6. χ^2 and p-values of the kinematic fitting at vertex A.

Decay mode	χ^2	p-value [%]
$^{16}\text{O} + \Xi^- \rightarrow {}^{10}_{\Lambda\Lambda}\text{Be} + {}^4\text{He} + t$	11.5	0.93
$^{16}\text{O} + \Xi^- \rightarrow {}^{11}_{\Lambda\Lambda}\text{Be} + {}^4\text{He} + d$	7.28	6.35
$^{16}\text{O} + \Xi^- \rightarrow {}^{12}_{\Lambda\Lambda}\text{Be} + {}^4\text{He} + p$	11.3	1.02

The probabilities of the three interpretations were evaluated with the chi-squared value of the kinematic fitting with a DOF of 3. The chi-squared and p-values for these three decay modes are summarized in Table 6. It is found that the most probable interpretation of this event is $^{11}_{\Lambda\Lambda}\text{Be}$ from the chi-squared values.

In the above analysis, it is assumed that no neutron was emitted at vertex C. The interpretation of vertex C is important because the analysis of vertex B is not effective in selecting possible candidates from the kinematics due to the large Q values of non-mesonic decays. If we assume neutron emission, the following decay modes are also accepted:

$${}^3_{\Lambda}\text{H} \rightarrow p + p + \pi^- + n, \quad (1)$$

$${}^4_{\Lambda}\text{H} \rightarrow d + p + \pi^- + n. \quad (2)$$

Table 7. Summary of $_{\Lambda\Lambda}$ Be double- Λ hypernuclei observed in past experiments. Multiple interpretations are listed for the MIKAGE and HIDA events.

Event	Target	Nuclide	$B_{\Lambda\Lambda}$ [MeV]	$\Delta B_{\Lambda\Lambda}$ [MeV]	Comment
MIKAGE [4]	^{12}C	$^6_{\Lambda\Lambda}\text{He}$	10.01 ± 1.71	3.77 ± 1.71	most probable
	^{12}C	$^{11}_{\Lambda\Lambda}\text{Be}$	22.15 ± 2.94	3.95 ± 3.00	
	^{14}N	$^{11}_{\Lambda\Lambda}\text{Be}$	23.05 ± 2.59	4.85 ± 2.63	
DEMACHIYANAGI [4]	^{12}C	$^{10}_{\Lambda\Lambda}\text{Be}^*$	11.90 ± 0.13	-1.52 ± 0.15	
HIDA [4]	^{14}N	$^{12}_{\Lambda\Lambda}\text{Be}$	22.48 ± 1.21		
	^{16}O	$^{11}_{\Lambda\Lambda}\text{Be}$	20.83 ± 1.27	2.61 ± 1.34	
Danysz [20–22]	^{12}C	$^{10}_{\Lambda\Lambda}\text{Be}$	14.7 ± 0.4	1.3 ± 0.4	

However, it should be noted that the branching ratio of $^3_{\Lambda}\text{H}$ and $^4_{\Lambda}\text{H}$ decay measured in a past experiment only show fewer than 30 instances of decay mode (1) and five of Eq. (2) among about 2000 $^3_{\Lambda}\text{H}$ and $^4_{\Lambda}\text{H}$ decays, respectively [17]. Moreover, a theoretical calculation also supports this small decay probability, e.g., 0.6% for decay mode (1) [18]. Thus, the $_{\Lambda\Lambda}\text{Li}$ nuclides in vertex A are very unlikely.

5. Discussion The newly observed double- Λ hypernuclear event is interpreted as the production of a $^{10}_{\Lambda\Lambda}\text{Be}$, $^{11}_{\Lambda\Lambda}\text{Be}$, or $^{12}_{\Lambda\Lambda}\text{Be}$ nucleus. The present result was compared with the results of past experiments. Candidates of $_{\Lambda\Lambda}\text{Be}$ double- Λ hypernuclei that were observed in past experiments are listed in Table 7.

The DEMACHIYANAGI event was interpreted as $^{10}_{\Lambda\Lambda}\text{Be}^*$ (most probable), an excited state of $^{10}_{\Lambda\Lambda}\text{Be}$, with a $B_{\Lambda\Lambda}$ value of 11.90 ± 0.13 MeV [4]. The present result of $^{10}_{\Lambda\Lambda}\text{Be}$ interpretation is consistent with the DEMACHIYANAGI event when considering that $^{10}_{\Lambda\Lambda}\text{Be}$ was generated in the ground state. The energy difference of the excited (2_1^+) and ground (0_1^+) states of the core nucleus ^8Be is 3.03 MeV [19], although both states are particle unbound. The double- Λ hypernucleus $^{10}_{\Lambda\Lambda}\text{Be}$ with $B_{\Lambda\Lambda} = 14.7 \pm 0.4$ MeV observed by Danysz is also consistent with the present result [20–22].

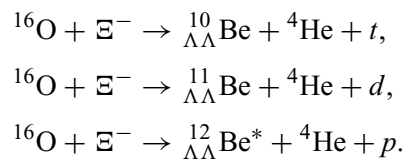
The present result of $^{11}_{\Lambda\Lambda}\text{Be}$ interpretation is also consistent with the past results given by HIDA and MIKAGE, in which $B_{\Lambda\Lambda}$ values were reported to be 22.15 ± 2.94 MeV, 23.05 ± 2.59 MeV, and 20.83 ± 1.27 MeV [4]. The present result has a small error on $B_{\Lambda\Lambda}$ because no neutron was emitted at the production vertex of the double- Λ hypernucleus.

In the case of $^{12}_{\Lambda\Lambda}\text{Be}$, the present analysis shows a negative $\Delta B_{\Lambda\Lambda}$ value (-2.7 ± 1.0 MeV), which is not consistent with NAGARA. However, if $^{12}_{\Lambda\Lambda}\text{Be}$ was produced in an excited state, $\Delta B_{\Lambda\Lambda}$ of the ground state is increased by its excitation energy. Therefore, the interpretation of $^{12}_{\Lambda\Lambda}\text{Be}$ cannot be excluded. The core nucleus, ^{10}Be , has excited states 3.368–6.263 MeV [19]. The level of $^{12}_{\Lambda\Lambda}\text{Be}$ is estimated to be similar to that of ^{10}Be (E. Hiyama, private communication). Thus, $\Delta B_{\Lambda\Lambda}$ for the ground state of $^{12}_{\Lambda\Lambda}\text{Be}$ becomes positive.

Three interpretations, $^{10}_{\Lambda\Lambda}\text{Be}$, $^{11}_{\Lambda\Lambda}\text{Be}$, and $^{12}_{\Lambda\Lambda}\text{Be}^*$, are accepted from the above consideration. Among them, $^{11}_{\Lambda\Lambda}\text{Be}$ is the most probable one from the analysis described in the previous section.

6. Summary An experiment to search for double- Λ hypernuclei with a counter-emulsion hybrid method, E07, was carried out at J-PARC. An impressive double- Λ hypernuclear event called “MINO” has been observed. Based on kinematic analysis, the nuclide of the double- Λ hypernucleus was

uniquely identified as a $\Lambda\Lambda$ Be. The event was interpreted as one of the following three candidates:



$B_{\Lambda\Lambda}$ ($\Delta B_{\Lambda\Lambda}$) of these double- Λ hypernuclei were obtained as 15.05 ± 0.11 MeV (1.63 ± 0.14 MeV), 19.07 ± 0.11 MeV (1.87 ± 0.37 MeV), and 13.68 ± 0.11 MeV (-2.7 ± 1.0 MeV), respectively, by assuming Ξ^- capture in the atomic 3D state with B_{Ξ^-} of 0.23 MeV. The negative $\Delta B_{\Lambda\Lambda}$ value of ${}^{12}_{\Lambda\Lambda}\text{Be}$ indicates that it was produced in the excited state. The most probable interpretation was found to be the production and decay of the ${}^{11}_{\Lambda\Lambda}\text{Be}$ nucleus from the kinematic fitting.

The emulsion scanning of the E07 experiment is ongoing. Twice as much statistics for Ξ^- stopping events than in E373 has been scanned and more than 10 events of double- and twin- Λ hypernuclei have been observed so far. Further impressive events are expected to be observed in the near future.

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