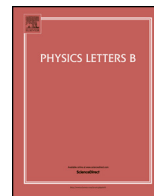




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Letter

Neutron radius determination of  $^{133}\text{Cs}$  and its impact on the interpretation of  $\text{CE}\nu\text{NS-CsI}$  measurement

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## ARTICLE INFO

Editor: H. Gao

## Keywords:

Elastic scattering  
Neutron radius  
Weak mixing angle

## ABSTRACT

Proton- $^{133}\text{Cs}$  elastic scattering at low momentum transfer is performed using an in-ring reaction technique at the Cooler Storage Ring at the Heavy Ion Research Facility in Lanzhou. Recoil protons from the elastic collisions between the internal  $\text{H}_2$ -gas target and the circulating  $^{133}\text{Cs}$  ions at 199.4 MeV/u are detected by a silicon-strip detector. The matter radius of  $^{133}\text{Cs}$  is deduced by describing the measured differential cross sections using the Glauber model. Employing the adopted proton distribution radius, a point-neutron radius of 4.86(21) fm for  $^{133}\text{Cs}$  is obtained. With the newly determined neutron radius, the weak mixing angle  $\sin^2\theta_W$  is independently extracted to be 0.227(28) by fitting the coherent elastic neutrino-nucleus scattering data. Our work limits the  $\sin^2\theta_W$  value in a range smaller than the ones proposed by the previous independent approaches, and would play an important role in searching new physics via the high precision  $\text{CE}\nu\text{NS-CsI}$  cross section data in the future.

The weak mixing angle,  $\sin^2\theta_W$ , is a fundamental parameter in the  $SU(2)_L \times U(1)_Y$  electroweak theory of the Standard Model (SM) [1]. Any deviation from the expected  $\sin^2\theta_W$  value in SM may serve as an indicative signature of new physics. Historically, the masses of top quark and Higgs boson were successfully predicted by the higher order diagram calculations with the measured  $\sin^2\theta_W$  [1].

To precisely constrain the  $\sin^2\theta_W$  value, various methods are developed to measure the dependence of  $\sin^2\theta_W$  on the transferred momentum [2,3]. Very recently, a high precision of about 0.1% for the  $\sin^2\theta_W$  determination was achieved by the CMS experiment [4]. At very low momentum transfer, for instance, one uses the atomic parity nonconservation (PNC) measurements of  $^{133}\text{Cs}$  [5,6]. However, as pointed out in Ref. [7], the  $\sin^2\theta_W$  determination via the PNC depends on many theoretical corrections. There is an obvious difference for the PNC amplitude correction associated with neutron skin effects [8]. The effects of the un-

known neutron distribution radius of  $^{133}\text{Cs}$  on the PNC were addressed about 25 years ago [9].

Recently the coherent elastic neutrino-nucleus scattering ( $\text{CE}\nu\text{NS}$ ) was observed using a  $\text{CsI}[\text{Na}]$  detector at the Oak Ridge National Laboratory [10]. The  $\text{CE}\nu\text{NS}$  measurement allows one to investigate many physics, for instance, neutrino nonstandard interactions [11], dark matter [12], and light vector  $Z'$  mediator [13], see Refs. [7,14] for details. Specially, the  $\text{CE}\nu\text{NS}$  experiment also provides a clear method to constrain  $\sin^2\theta_W$  at low momentum transfer [11]. Different from other experimental methods [2,3], the  $\sin^2\theta_W$  deviation between the  $\text{CE}\nu\text{NS}$  experiment and the SM prediction could give a hint on new physics, e.g.  $\nu$ -nucleon interactions [15]. Although majority of observables were experimentally measured, there was still a lack of precise neutron radii of  $^{133}\text{Cs}$  and  $^{127}\text{I}$  [11,13,16] in interpreting the  $\text{CE}\nu\text{NS-CsI}$  data. As underscored by the COHERENT collaboration [11], the uncertainty in the  $\text{CE}\nu\text{NS-CsI}$  cross section calculations is dominated by the neutron dis-

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Received 7 April 2024; Received in revised form 14 July 2024; Accepted 21 July 2024

Available online 26 July 2024

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tributions of  $^{133}\text{Cs}$  and  $^{127}\text{I}$ . The influence of the neutron radii on the CEvNS-CsI cross section data interpretation was also addressed by other independent investigations [16]. In a word, inaccurate or biased treatment of the neutron radii of  $^{133}\text{Cs}$  and  $^{127}\text{I}$  would lead to the misidentification of possible signals of new physics [16]. Usually,  $\sin^2\theta_W$  can be deduced through a fit to the CEvNS-CsI data by fixing neutron radii from theoretical predictions of various models [11,17–19]. Alternatively, the neutron radii can be deduced by assuming a fixed  $\sin^2\theta_W$  value [20]. Up to date, the reported (average) neutron radii of  $^{133}\text{Cs}$  spread from 4.6 fm through 6.6 fm [15,17,18,20–27].

The accurate determination of the neutron radius for  $^{133}\text{Cs}$  has garnered significant attention across the atomic, nuclear, and particle physics communities owing to its significance in fundamental researches. However, it is challenging to experimentally measure the neutron radius of  $^{133}\text{Cs}$  in normal kinematics, due to the low melting point of 28 °C and spontaneous ignition in air. In this Letter, we introduce an innovative approach to determine the neutron radius of  $^{133}\text{Cs}$  by measuring the proton elastic scattering at low momentum transfer using an in-ring reaction technique and inverse kinematics. We elucidate the impact of this determination on the  $\sin^2\theta_W$  extraction from the recent CEvNS-CsI data.

The novel in-ring reaction experiment was carried out at the experimental Cooler Storage Ring (CSR) of the Heavy Ion Research Facility in Lanzhou (HIRFL) [28]. Such kinds of experiments are especially suited for the small-angle differential cross section measurements at low momentum transfer [29,30]. CSR, which is equipped with the electron cooler [31] and the internal  $\text{H}_2$ -gas-jet target [32], was operated for an in-ring reaction experiment at a magnetic rigidity of about 5.205 Tm. The  $^{133}\text{Cs}^{27+}$  beam with an energy of 204 MeV/u from the main storage ring (CSRm) was stripped off all bound electrons utilizing an aluminum foil with a thickness of 0.21 mm at the radioactive ion beam line (RIBLL2). Then the  $^{133}\text{Cs}^{55+}$  ions with an energy of about 199.4 MeV/u were transported through RIBLL2, and injected into CSR. The stored  $^{133}\text{Cs}^{55+}$  ions in CSR interacted repeatedly with the  $\text{H}_2$ -gas target of about  $10^{12}$  atoms/cm<sup>2</sup> thickness [32]. The electron cooling at CSR was operated to compensate for the energy loss of the ions caused by the collisions with the gas target and residual gas. The recoil protons from the proton- $^{133}\text{Cs}$  elastic scattering within angular range from about 85° to 90° in the laboratory system were measured by a double-sided silicon-strip detector (DSSD) with a typical energy resolution of better than 1%. The employed 1000  $\mu\text{m}$  thick DSSD had an active area of  $64 \times 64$  mm<sup>2</sup> and was segmented into  $32 \times 32$  strips. The proton energy and detection efficiency were calibrated by radioactive sources [30]. Fig. 1 illustrates the scatter plot of the recoil proton energy versus the strip number of DSSD. Further experimental details were described in our previous works [30,33,34].

Given the fact that the flight paths and energies of recoil protons are hardly altered by secondary collisions with the thin gas target, the relative small-angle differential cross sections  $\frac{d\sigma}{d\Omega}(\theta)$  of proton- $^{133}\text{Cs}$  elastic scattering are determined via [30]

$$\frac{d\sigma}{d\Omega}(\theta) = \frac{1}{\sin\theta} \left( \frac{\Delta N_{\text{all}}}{\Delta\theta} - \frac{\Delta N_{\text{bg}}}{\Delta\theta} \right), \quad (1)$$

where  $\Delta N_{\text{all}}$  is the number of all measured events in the scattering angle interval  $\Delta\theta$  in the center-of-mass (c.m.) frame, and  $\Delta N_{\text{bg}}$  is the corresponding background estimated by the measured events in the rectangles enclosed by the yellow solid lines, see Fig. 1. The scattering angles  $\theta$  were determined by the proton kinematic energies  $K_{\text{lab}}$  via the relation of  $2m_p K_{\text{lab}} = 2P^2(1 - \cos\theta)$ , with  $m_p$  and  $P$  being the proton rest mass and the c.m. momentum, respectively. The uncertainties of  $\theta$  are smaller than the used  $\Delta\theta$  value of 0.1°. To reduce the effects of solid angle and detection efficiency [30], only single coincidence events between  $X$  and  $Y$  strips within the energy range of about 0.6 MeV through 3.6 MeV were considered. As shown in Fig. 1, the events of elastic scattering are mixed, to some extent, with those of inelastic scattering associated with the low-lying 0.08- and 0.16-MeV excited states [35]. However, experi-

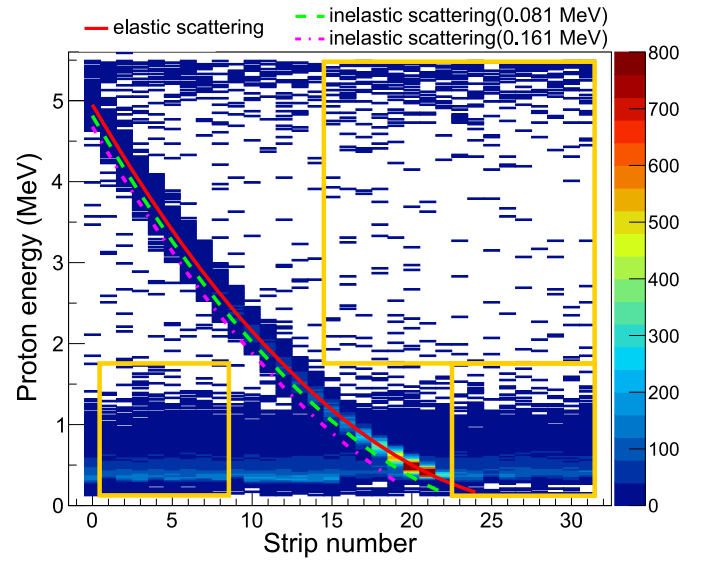


Fig. 1. Scatter plot of the recoil proton energy versus the strip number of DSSD. The solid (red), dashed (green), and dash-dotted (pink) lines denote the calculated proton energies for elastic and two inelastic scattering channels, respectively. For more details see text.

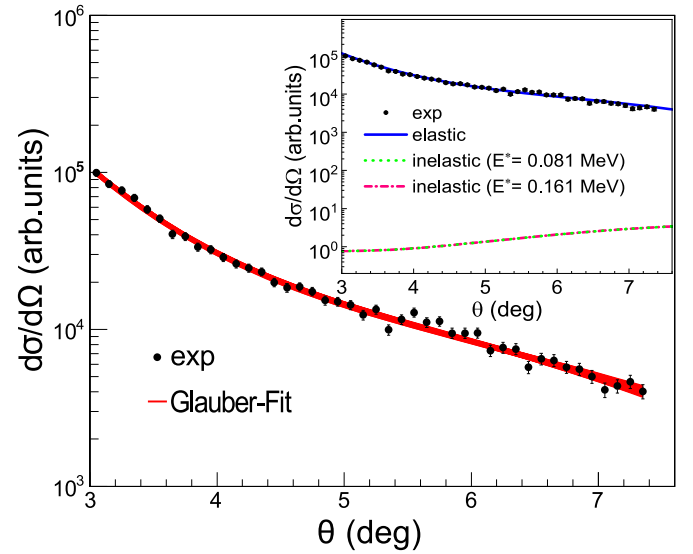


Fig. 2. The measured  $\frac{d\sigma}{d\Omega}(\theta)$  for proton- $^{133}\text{Cs}$  elastic scattering and the  $2\sigma$  distribution of fit curves (red). The insert shows the calculated elastic and inelastic differential cross sections using FRESKO [38]. The trend of the measured  $\frac{d\sigma}{d\Omega}(\theta)$  values agrees well with the FRESKO calculations. Conveniently  $\frac{d\sigma}{d\Omega}(\theta)$  are normalized to the FRESKO calculations.

ments have already shown that the inelastic scattering cross sections are very tiny in the small angular region, compared to the elastic scattering cross sections [36,37]. For the  $^{133}\text{Cs}$  case, according to the FRESKO calculations [38] with the phenomenological optical model [39], the inelastic scattering cross sections are also several orders of magnitude smaller than those of the elastic scattering in the measured small angular range, see the insert of Fig. 2. Therefore, the contribution of the inelastic scattering can be safely ignored, compared to the several percent relative uncertainties of measured differential cross sections. The measured small-angle  $\frac{d\sigma}{d\Omega}(\theta)$  values are shown in Fig. 2.

It is well known that the small-angle elastic  $\frac{d\sigma}{d\Omega}(\theta)$  distributions are sensitive to matter distribution radius [29,30,40–43]. Especially, the reaction mechanism is relatively simple at small angles and thus cor-

rection terms of reaction models are negligible compared to the cases of large angle scattering [44,45]. Therefore, many experimental setups at various facilities are developed to determine the matter radii by measuring the small-angle  $\frac{d\sigma}{d\Omega}(\theta)$  of proton-nucleus elastic scattering [29,30,33,40–42]. In the present work, a well established procedure [44,46,47] based on the Glauber multiple-scattering theory [48] is employed to extract the matter radius of  $^{133}\text{Cs}$  through describing the measured  $\frac{d\sigma}{d\Omega}(\theta)$ . The  $\frac{d\sigma}{d\Omega}(\theta)$  values are expressed in the Glauber model as a function of the matter density distribution  $\rho(r)$  and the proton-nucleon scattering amplitude  $f_{pi}(q)$  with  $i = n$  or  $p$ , see Ref. [47] for details. To address the model-dependent errors related to the hadronic probes [49] as much as possible, similar to these works [30,46,50,51], the scattering amplitude parameters of the Glauber model were calibrated at 200 MeV [51] to be  $\sigma_{pp} = 1.788(20) \text{ fm}^2$ ,  $\sigma_{pn} = 3.099(27) \text{ fm}^2$ ,  $\alpha_{pp} = 0.893(17)$ ,  $\alpha_{pn} = 0.325(23)$ , and  $\beta_{pp} = \beta_{pn} = 0.528(41) \text{ fm}^2$ , which are adopted here to calculate the  $f_{pi}(q)$ . Due to the intermediate energy, similar to previous works [30,51], the effects of Pauli blocking on  $\sigma_{pp}$  and  $\sigma_{pn}$  are corrected. These values have been adopted to fit the differential cross sections of proton- $^{16}\text{O}$  elastic scattering at 200 MeV, and the resulting matter radius of  $^{16}\text{O}$  is comparable with other experimental results [51].

In the radius fitting procedure, the  $\rho(r)$  is described by the two-parameter Fermi model as

$$\rho(r) = \frac{\rho(0)}{1 + \exp[(r - R)/a]}, \quad (2)$$

with  $\rho(0)$ ,  $R$ , and  $a$  being the density normalization factor, half-density radius, and diffuseness parameter, respectively. Although the  $R$  and  $a$  values cannot be simultaneously constrained by  $\frac{d\sigma}{d\Omega}(\theta)$  in small angular range [30], the matter radius is almost independent of  $a$  in the range of 0.50 fm through 0.55 fm for the medium-heavy nuclei [47]. Thus, following the methods in [30,47,52,53], we set  $a = 0.53(3) \text{ fm}$ , which was deduced from the neighboring  $^{116,124}\text{Sn}$  and  $^{208}\text{Pb}$  nuclei [47]. Additionally, it is challenging to accurately determine reaction luminosity for the in-ring reaction experiments, not only due to the changes of the gas-target density and beam intensity in time but also due to the uncertainty on the overlap between beam and gas target [54]. Consequently, the theoretical X-ray atomic cross sections were used to deduce the in-ring reaction luminosity from time to time [33,55,56]. However, an unreasonable luminosity would result in a radius deviation [47]. Therefore, a free cross section normalization factor  $L_0$  is introduced as in Refs. [30,43]. Benefiting from such treatment, the radius uncertainty from the unreasonable absolute  $\frac{d\sigma}{d\Omega}(\theta)$  normalization can be reduced [47]. Subsequently,  $R$  and  $L_0$  are freely adjusted to fit the experimental  $\frac{d\sigma}{d\Omega}(\theta)$  with the Glauber model, and  $\rho(0)$  is also accordingly varied to keep the mass number  $A=133$ .

As shown in Fig. 2, the measured  $\frac{d\sigma}{d\Omega}(\theta)$  are well described with the Glauber model by adjusting  $R$  and  $L_0$ . With the obtained  $R$  and fixed  $a$ , a root-mean-square (rms) point-matter radius  $R_{\text{pm}}$  for  $^{133}\text{Cs}$  is determined to be

$$R_{\text{pm}} = \left( \frac{\int \rho(r)r^4 dr}{\int \rho(r)r^2 dr} \right)^{\frac{1}{2}} = 4.811 \pm 0.127 \text{ fm}, \quad (3)$$

where uncertainties from statistics, input parameters, and Glauber model are about 0.12 fm, 0.03 fm, and 0.03 fm, respectively. The radius uncertainty caused by statistics is the standard deviation of radius distribution with a confidence level of 95.5%, which is obtained by fitting the randomly sampled experimental  $\frac{d\sigma}{d\Omega}(\theta)$  within  $2\sigma$  band based on the Gaussian distribution [47], while other input parameters are unchanged. The uncertainty resulted by input parameters (namely,  $\sigma$ ,  $\alpha$ ,  $\beta$ , and  $a$ ) is also deduced with the similar method, but the radius distribution is obtained by fitting the unchanged experimental  $\frac{d\sigma}{d\Omega}(\theta)$  with the randomly sampled input parameters within corresponding  $2\sigma$  band. In principle, the model-dependent error can be estimated by comparing the measured results with the standard true values. However, due to

lack of well-known matter radius data, it is challenging to estimate the model-dependent error at present. In this work, the model-dependent error at 200 MeV is estimated by comparing the well-known proton radii with the matter radii of  $^{12}\text{C}$ ,  $^{16}\text{O}$ , and  $^{28}\text{Si}$  determined with the similar method, where similar proton and matter radii are expected for the  $N = Z$  nuclei. To check the effects of background, only recoil protons with energies  $> 1 \text{ MeV}$  were analyzed, and a consistent radius of 4.825 fm is obtained. Details and reliability considerations about radius determinations can be found in Refs. [30,47].

With the obtained  $R_{\text{pm}}$ , a point-neutron distribution radius  $R_{\text{pn}}$  of  $^{133}\text{Cs}$  is determined to be

$$R_{\text{pn}} = \sqrt{\frac{A}{N} R_{\text{pm}}^2 - \frac{Z}{N} R_{\text{pp}}^2} = 4.86 \pm 0.21 \text{ fm}, \quad (4)$$

where  $N$ ,  $Z$ , and  $A$  are the neutron, proton, and mass number, respectively. The adopted point-proton radius  $R_{\text{pp}}$  of 4.740(5) fm for  $^{133}\text{Cs}$  is deduced from charge radius [30,57].

We extract the neutron skin of  $^{133}\text{Cs}$  to be  $R_{\text{pn}} - R_{\text{pp}} = 0.12(21) \text{ fm}$ . Meanwhile, using the neutron skin linear relationship between  $^{133}\text{Cs}$  and  $^{208}\text{Pb}$ , established by various effective interactions [58], we also derive the weight-average neutron skin of  $^{133}\text{Cs}$  to be 0.14(3) fm from the Parity-violating electron scattering results of  $^{48}\text{Ca}$  and  $^{208}\text{Pb}$  [59,60], where the data for  $^{48}\text{Ca}$  [60] is related to the neutron skin of  $^{208}\text{Pb}$  using the relationship shown in Fig. 5 in Ref. [60]. The two values are consistent with each other, and agree to the value of 0.13(4) fm calculated by the empirical linear relationship from the antiprotonic atom experiment [61]. The latter was adopted to estimate the PNC amplitude correction associated with the neutron skin effects [62].

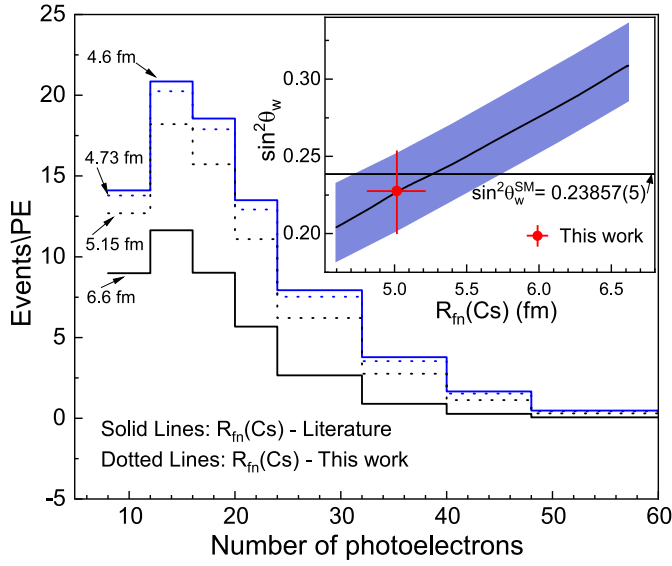
Compared to the SM prediction, any deviation of the effective  $\sin^2\theta_W$  determined by the CEvNS-CsI data would be an indicative signature of new physics associated with, e.g. the  $\nu$  interactions. Now we discuss the impact of the neutron radius on the  $\sin^2\theta_W$  determination employing the CEvNS-CsI cross section data [11]. For the description of the CEvNS-CsI data in SM,  $\sin^2\theta_W$  and the neutron distribution form factor  $F_n(q^2)$  are involved, see Eq. (1) in Ref. [27] for details. The folded-neutron radii  $R_{\text{fn}}(\text{Cs})$  and  $R_{\text{fn}}(\text{I})$  for  $^{133}\text{Cs}$  and  $^{127}\text{I}$  are indispensable to determine the corresponding  $F_n(q^2)$ . Previous analyses indicated that different parameterizations of nuclear form factors would not lead to different  $\sin^2\theta_W$  [13,15,27]. In this work, the Helm form factor is adopted to calculate the expected CEvNS-CsI signal event number, which is related to  $R_{\text{fn}}$  through the diffraction radius  $R_0$  defined as [27]

$$R_0(\text{K}) = \sqrt{5/3[R_{\text{fn}}^2(\text{K}) - 3s^2]}, \quad (5)$$

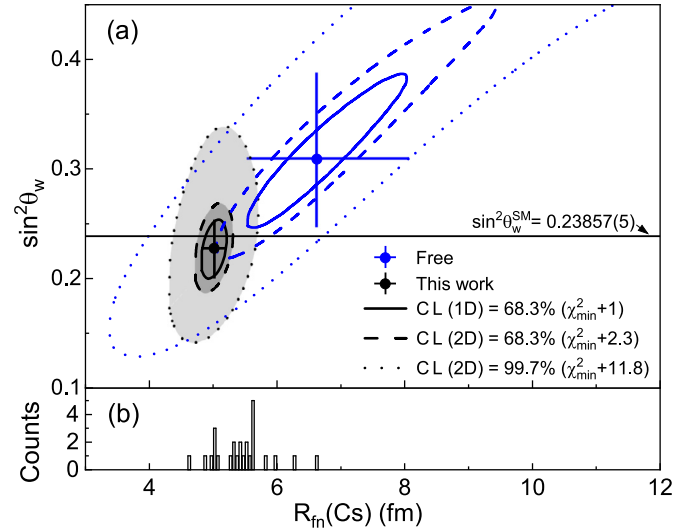
where  $s = 0.9 \text{ fm}$  is the surface thickness and  $\text{K} = \text{I}$  or  $\text{Cs}$ . The  $R_{\text{fn}}$  can be directly determined from the point-neutron radius via  $[R_{\text{pn}}^2 + (0.864 \text{ fm}^2)]^{1/2}$  [19].

Fig. 3 depicts the CEvNS-CsI signal event number as a function of the photoelectron number expected with different  $R_{\text{fn}}(\text{Cs})$ . In this work, the radius  $R_{\text{fn}}(\text{I})$  is determined using the  $R_{\text{fn}}(\text{Cs})$  value via  $R_{\text{fn}}(\text{Cs}) - R_{\text{fp}}(\text{Cs}) + R_{\text{fp}}(\text{I})$  considering that  $^{133}\text{Cs}$  and  $^{127}\text{I}$  have almost the same neutron skin thicknesses because of similar  $(N - Z)/A$  values [61], where the folded-proton radius  $R_{\text{fp}}(\text{K})$  is determined by the well-known charge radius [19,57]. It is evident that the expected event spectrum is significantly affected by the adopted neutron radius of  $^{133}\text{Cs}$ , which is similar as addressed in Ref. [16]. As indicated in the inset of Fig. 3, there exists a strong correlation between the values of  $R_{\text{fn}}(\text{Cs})$  and  $\sin^2\theta_W$  in the CEvNS-CsI data analysis. An incorrect neutron radius can thus introduce significant shift in the estimation of  $\sin^2\theta_W$ .

Previously, a method was used combining the CEvNS and PNC data to precisely constrain  $\sin^2\theta_W$  [27]. However, possible new physics in either the CEvNS or PNC process may be ignored in such analysis due to the assumption that the two processes give the same  $\sin^2\theta_W$  value. In addition, the PNC amplitude corrections depend on theory [7], which may result in different values of  $\sin^2\theta_W$  [27].



**Fig. 3.** The expected CEvNS-CsI signal event number as a function of the photoelectron number calculated by using different neutron radii of  $^{133}\text{Cs}$ . The blue and black solid lines are obtained by using the reported (average) neutron radius values of 4.6 fm [18] and 6.6 fm [27] for  $^{133}\text{Cs}$ , respectively, see Fig. 4(b). The blue and black dotted lines are calculated with the presently deduced radii  $R_{\text{fn}}^{\text{exp}}(\text{Cs})$  of  $4.94 \pm 0.21$  fm, respectively. The insert shows a correlation between  $R_{\text{fn}}(\text{Cs})$  and  $\sin^2\theta_W$ , where the blue area is the uncertainty caused by the CEvNS-CsI data [11]. The red symbol shows the obtained  $R_{\text{fn}}^{\text{fit}}(\text{Cs})$  and  $\sin^2\theta_W$  by Eq. (6), details see text.



**Fig. 4.** (a) The  $\chi_{\text{all}}^2$  contours in the plane of  $R_{\text{fn}}$  versus  $\sin^2\theta_W$ . The blue curves and point represent results when both  $R_{\text{fn}}$  and  $\sin^2\theta_W$  are free variables in the CEvNS-CsI data fitting. The black curves and point add the constraint imposed by the presently deduced radius. (b) The distribution of reported neutron radii of  $^{133}\text{Cs}$  [15,17,18,20–27] deduced from the CEvNS-CsI data [10,11].

To avoid any improper input of neutron radius and  $\sin^2\theta_W$ , we performed an independent two-dimensional (2D) fit as Refs. [15,27] using both  $R_{\text{fn}}$  and  $\sin^2\theta_W$  as free adjustable parameters. The  $\sin^2\theta_W$  value from the best fit is shown in Fig. 4 as the blue symbol. The blue curves represent the distributions of the  $\sin^2\theta_W$  values obtained from fitting procedures under different confidence levels (CL). Because of the strong correlation of the two free parameters, the independent 2D fit can not yield a well-constrained  $\sin^2\theta_W$  value, nor the realistic neutron radius, as shown in Fig. 4. Therefore accurate neutron radius is essential to deduce the  $\sin^2\theta_W$  value and then to search for new physics beyond SM.

To show the impact of our newly determined neutron radius on the  $\sin^2\theta_W$  value, the expected CEvNS events are fitted to the experimental events by adjusting  $\sin^2\theta_W$  and  $R_{\text{fn}}^{\text{fit}}(\text{Cs})$  via the chi-square function defined here as

$$\chi_{\text{all}}^2 = \chi_{\text{Csl}}^2 (R_{\text{fn}}^{\text{fit}}(\text{Cs}), \sin^2\theta_W) + \left( \frac{R_{\text{fn}}^{\text{fit}}(\text{Cs}) - R_{\text{fn}}^{\text{exp}}(\text{Cs})}{\sigma^{\text{exp}}} \right)^2, \quad (6)$$

where  $\sigma^{\text{exp}}$  is the uncertainty of  $R_{\text{fn}}^{\text{exp}}(\text{Cs})$ . The  $\chi_{\text{Csl}}^2$  was constructed (see Ref. [27] for details) and used for the free two-dimensional fit as mentioned above. The used  $R_{\text{fn}}^{\text{exp}}(\text{Cs})$  of 4.936(210) fm is obtained by the presently deduced point-neutron radius. Now, for the first time we fix the range of neutron radii in the second term of Eq. (6). As a consequence,  $\sin^2\theta_W$  of 0.227(28) is independently extracted from the best fit to the CEvNS-CsI data [11] as shown with the black symbol in Fig. 4. The black curves represent the distributions of the  $\sin^2\theta_W$  values obtained from fitting procedures for different CL. Our determined value  $\sin^2\theta_W = 0.227(28)$  agrees within error bars with the SM prediction of 0.23857(5) at low momentum transfer [63]. Compared to the  $\sin^2\theta_W$  value of  $0.309_{-0.063}^{+0.078}$  determined by the free independent 2D fit, our analysis well constrains the  $\sin^2\theta_W$  value and improves its precision by a factor of about 2.5. However, the present precision of about 10% for our  $\sin^2\theta_W$  value is still too large to effectively probe new physics. This large uncertainty is mainly caused by the CEvNS-CsI data [11], see the red symbol and blue area in the inset of Fig. 3. Therefore, the key neutron radius value provided in this work would play an important role in searching new physics related to the  $\nu$  interactions via the high precision CEvNS-CsI data in the future.

In conclusion, the proton elastic scattering off  $^{133}\text{Cs}$  at 199.4 MeV/u was investigated in inverse kinematics at HIRFL-CSR. Combined with the proton distribution radius, a point-neutron radius of 4.86(21) fm for  $^{133}\text{Cs}$  was extracted. For the first time we fix the range of neutron radii of  $^{133}\text{Cs}$  in the CEvNS-CsI data analysis, and consequently the weak mixing angle was extracted independently with higher accuracy than the ones proposed by the previous independent approaches. This work provides a key neutron radius data and constitutes a new approach for the interpretation of high precision CEvNS-CsI data in the future.

#### Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

#### Data availability

Data will be made available on request.

#### Acknowledgements

We thank the staffs in the accelerator division for providing a stable beam. This work is supported in part by the National Key R&D Program of China (Grant No. 2023YFA1606401), by the NSFC (Grants No. 12375115, No. 12022504, No. 12121005, No. 12075255, and No. 12275186), and by the CAS Open Research Project of large research infrastructures.

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