



Full Length Article

The upgraded summing NaI(Tl) (SuN++) absorption spectrometer

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ABSTRACT

Simulations of astrophysical processes require a plethora of nuclear physics input. In particular, models of neutron-capture nucleosynthesis like the *s*, *i*, and *r* processes require β -decay information and experimentally constrained neutron-capture reaction rates. Past experiments with the 4π Summing NaI(Tl) (SuN) total absorption spectrometer have provided these physics quantities. Here, we outline an upgrade of SuN to SuN++, where 20 new segments (12 NaI(Tl) and 8 CeBr₃) have been integrated into the pre-existing SuN total absorption spectrometer to provide increased energy and time resolution in β -decay experiments. The details of the newly upgraded SuN++ total absorption spectrometer are discussed with results from the commissioning experiment at the Facility for Rare Isotope Beams (FRIB) utilizing a ⁷⁰Cu beam.

1. Introduction

The astrophysical origin of over half of the elements in the universe heavier than iron is largely attributed to neutron-capture nucleosynthesis processes. Nuclear data play a significant role in the calculation of elemental abundances produced by these astrophysical nucleosynthesis processes, and in the absence of experimental measurements,

theoretical calculations of nuclear physics quantities are used as model input. Theoretical values often lead to large uncertainties in calculated abundances, notably for the *r* [1] and *i* [2,3] processes, which occur farther from the valley of stability [4–6]. Current astrophysical models of *i*-process nucleosynthesis rely heavily on neutron-capture (*n*, γ) reactions, and the dominating source of uncertainty arising

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from nuclear physics input in *i*-process nucleosynthesis simulations comes from theoretical predictions of these (n, γ) rates on unstable neutron-rich nuclei [6,7]. Multiple sensitivity studies have indicated that quantities such as nuclear masses, neutron-capture reaction rates, β -delayed fission, and half-lives all play integral roles in simulations of neutron-capture and β -decay competition in *r*-process nucleosynthesis pathways [4,5,8]. Uncertainties in *r*-process abundance patterns from modern simulations persist even when all trajectories for a simulation are combined, further emphasizing the importance of incorporating nuclear data into these simulations [9].

Simulations of the *r* process can be additionally impacted by experimental measurements of β -decay properties, including β -delayed neutron-emission probabilities, half-lives, and β -feeding values [10–15]. In the cases where experimental information is unavailable, which is common for the neutron-rich isotopes involved in the *r* process, theoretical calculations are used to predict these β -decay properties [16–19]. Theoretical calculations of β -decay properties are typically benchmarked against known half-lives and β -delayed neutron emission probabilities, which do not fully reflect the β -decay strength of neutron-rich nuclei. A more-direct approach is to compare experimental β -feeding intensity (I_β) values with theory.

I_β values are typically determined from high-resolution γ -ray measurements with high-purity germanium (HPGe) detectors. Such detectors have a low intrinsic efficiency, resulting in weak branching measurements to high-lying states in the child nucleus, causing inflated β -feeding to lower-lying states. This phenomenon is commonly referred to as the Pandemonium effect [20]. More accurate I_β values can be determined by total absorption spectroscopy (TAS), where a large volume, high efficiency detector is used to measure γ rays following β decays [21,22].

TAS analysis can be performed using a segmented total absorption spectrometer, such as Summing NaI(Tl) total absorption spectrometer (SuN) [23]. The experimental capabilities of SuN to disentangle β -decay information from neutron-rich nuclei far from stability, like those produced at the Facility for Rare Isotope Beams (FRIB), is limited by its current segmentation and resolution. Here, we present the details of an upgrade from SuN to the SuN++ total absorption spectrometer, which aims to address these problems. We present the characterization of the upgraded SuN++ calorimeter and details of the commissioning run at FRIB with a beam of ^{70}Cu .

2. The upgraded summing NaI(Tl) total absorption spectrometer (SuN++)

2.1. The SuN total absorption spectrometer

The SuN total absorption spectrometer is a 16 inch by 16 inch cylindrical total absorption spectrometer comprised of eight optically isolated NaI(Tl) segments. SuN was originally developed with the purpose of measuring capture reactions on relevant nuclei in astrophysical reactions. Details of the characterization and commissioning of SuN were reported by Simon et al. in Ref. [23].

Since its commissioning, SuN has been used to perform total absorption spectroscopy on multiple *r*-process nuclei [24–26] and to constrain several *i*- and *r*-process neutron-capture reaction rates with the β -Oslo Method [27–32]. The resolution of SuN (6.1(2)% at 662 keV) [23] significantly limits the ability to build new level schemes when measuring the β decay of exotic nuclei. The limited number of segments in SuN (eight) can also make it difficult to disentangle high-multiplicity cascades. Furthermore, the NaI(Tl) crystals in SuN are sensitive to electrons from β decay, and thus the β -decay electron-induced background detected in SuN can further complicate studies of neutron-rich β -decaying nuclei. In previous and current studies with SuN and SuN++, β -decay electrons have been included in Geant4 simulations of the detectors, and considered in the analysis process.

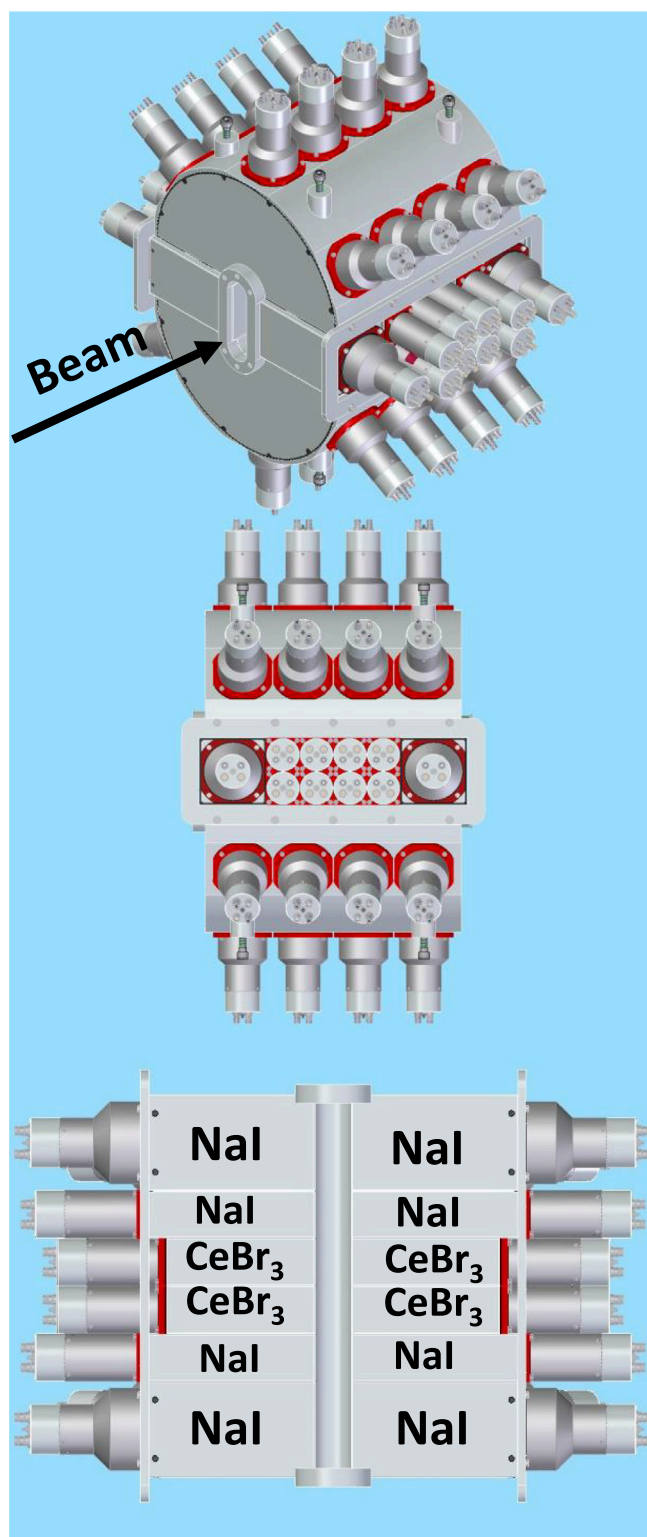


Fig. 1. Schematic drawing of the SuN++ detector (a), a side-view (b), and a center-view (c) where the new detector additions can be seen.

2.2. The SuN++ total absorption spectrometer

To resolve the aforementioned problems, 20 additional segments (12 NaI(Tl) segments and 8 CeBr₃ segments) were added to the original SuN total absorption spectrometer. The CeBr₃ segments are meant to improve the timing and energy resolution of the TAS detector to help

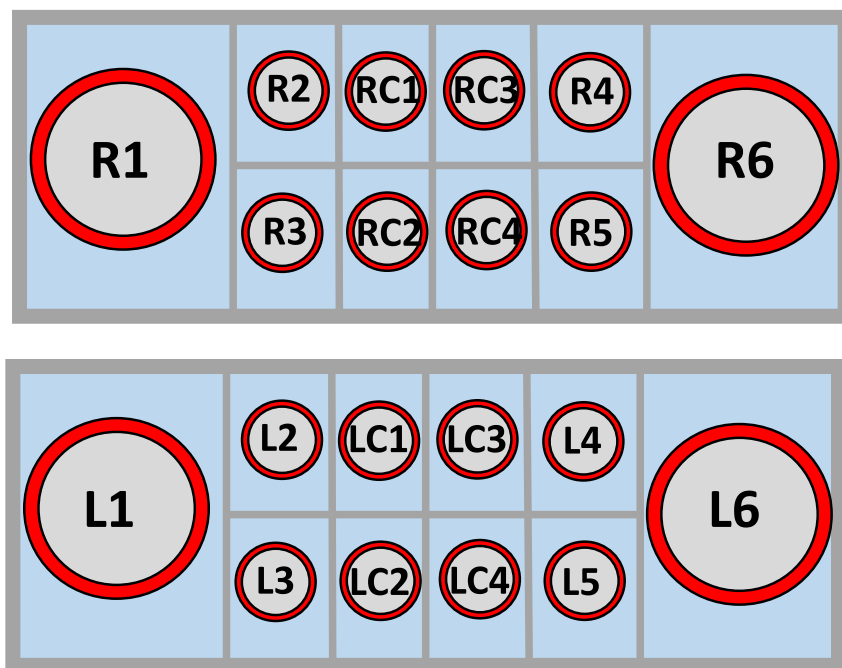


Fig. 2. Naming schematic of the SuN++ segments. The detectors on the right side of SuN++ are referred to as “R” detectors and on the left as “L” detectors. A “C” is added before the CeBr₃ detector to differentiate them the NaI(Tl) detectors.

build level schemes. With the addition of 12 NaI(Tl) segments, the 20 total new segments provide a higher level of segmentation that can be used to differentiate high-multiplicity cascades. Future plans to add electron detectors to the SuN++ system will allow for the capability to track and reject β -decay electrons in the total absorption spectrometer.

The SuN++ upgrade was manufactured by SCIONIX and was developed at FRIB. The eight CeBr₃ segments are 152 mm long by 51 mm wide square prisms, enclosed in stretched polytetrafluoroethylene (PTFE) and covered in a 1 mm thick aluminum casing. At the optical interface, a quartz window is optically coupled to a 51 mm diameter Hamamatsu PMT (Model No. R6231) with a head-on type connection. Two different sizes of NaI(Tl) segments were also included in SuN++, four “large” NaI(Tl) and eight “small” NaI(Tl) segments. The large NaI(Tl) segments are 178 mm long by 102 mm wide square prisms, while the small NaI(Tl) segments are 178 mm long by 51 mm wide square prisms. Each of these detector crystals is also enclosed in stretched PTFE and covered in 1 mm thick aluminum casing. The quartz window in the “small” NaI(Tl) segments is optically coupled to the Model No. R6231 Hamamatsu PMT. The quartz window in the “large” NaI(Tl) segments is optically coupled to a 90 mm diameter Hamamatsu PMT (Model No. R14689), also with a head-on type connection.

On each side of SuN++ relative to the beam direction, ten new segments are inserted between the existing halves of SuN (two large NaI(Tl), four small NaI(Tl), and four CeBr₃ detectors), seen in the schematic of SuN++ in Fig. 1. Two large NaI(Tl) segments are located at both ends of the detector. Two small NaI(Tl) segments are stacked upon each other and are placed on the innermost edge of the two large NaI(Tl) segments, followed by four CeBr₃ segments that are positioned at the middle of SuN++. This layout is mirrored on the other half of the detector. The naming-scheme for the SuN++ segments is illustrated in Fig. 2.

A schematic of the SuN++ electronics and data acquisition system is shown in Fig. 3. Each of the segments from the original SuN detector is read by three photomultiplier tubes (PMTs) positioned around the segment, giving a total of 24 optically coupled PMTs. All of the new SuN++ segments have individual PMTs that are optically coupled to the crystals. A WIENER MPOD [33] high voltage power supply is used to provide bias voltage to all 44 PMTs. The 24 PMT signals from SuN are

sent to a preamplifier and then into the FRIB Digital Data Acquisition System (DDAS) [34]. The 20 photomultiplier signals from SuN++ are fed directly into DDAS. A coincidence condition is set between the three PMTs that comprise each optically isolated segments of the original SuN detector in the digitizer firmware. The SuN++ segments were all high-voltage gain-matched such that the centroids of the 1460 keV ⁴⁰K peak (from room background) were at the same channel number.

3. Results

3.1. Summing efficiency and source tests

¹³⁷Cs and ⁶⁰Co γ -sources were used to test the resolution of the individual segments and the total detection efficiency of the SuN++ total absorption spectrometer. The ¹³⁷Cs source emits a γ ray at 662 keV, and ⁶⁰Co emits two γ rays at 1173 and 1332 keV. For the new individual segments of SuN++, the energy resolution values are reported in Fig. 4. The average resolution for the 1332 keV single- γ line from ⁶⁰Co is 4.65(8)% for the large NaI(Tl), 4.71(16)% for the small NaI(Tl), and 3.38(4)% for the CeBr₃ detectors, compared to 5.8(2)% for SuN [23]. To characterize SuN++ and compare its performance to expectations, the Geant4 simulation toolkit was used [35]. This Geant4 simulation was built upon the SuN simulation for β -decay experiments outlined by Dombos et al. [24].

The ⁶⁰Co source was also used to determine the efficiency of detecting a single γ ray at an energy of 1173 keV at the center of SuN++. The single γ -efficiency was experimentally determined to be 71.2(8)%, which agrees within error with the single γ -efficiency determined from Geant4 (70.2(5)%). A typical ⁶⁰Co total absorption spectrum, sum of segments spectrum, and multiplicity spectrum are shown in Fig. 5(a), (b), and (c), respectively. The total absorption spectrum is made by summing all energy deposited in SuN++ within an event window; this spectrum is related to the excitation energy populated in the child nucleus. The sum-of-segments spectrum is comprised of the sum of histograms of the SuN++ segments (related to the individual γ -rays that de-excite from the child nucleus) which is made by adding the counts in the spectrum from each individual segment together. The multiplicity spectrum shows the number of SuN++ segments that register a signal

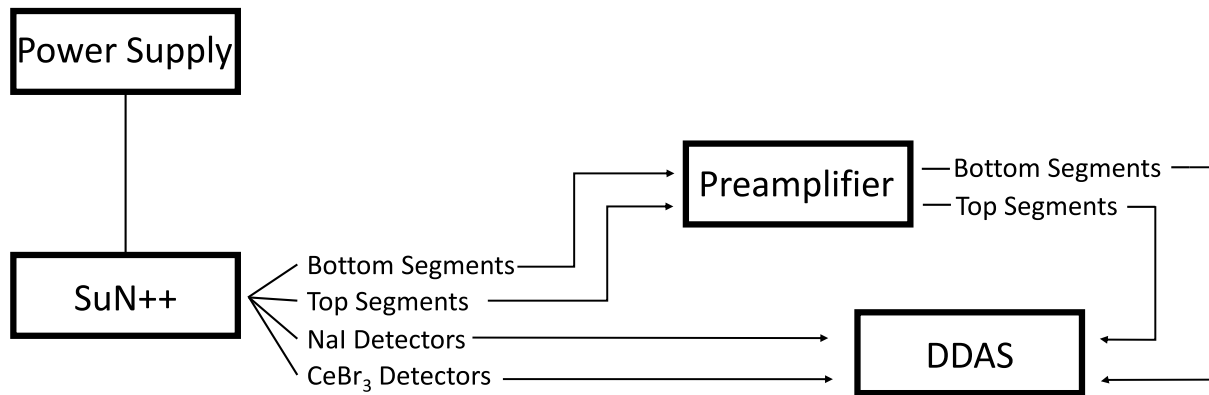


Fig. 3. Schematic of the SuN++ electronics. Signals from the original SuN PMTs are sent through a preamplifier before being sent to DDAS. Signals from the new SuN++ detectors are fed directly to DDAS.

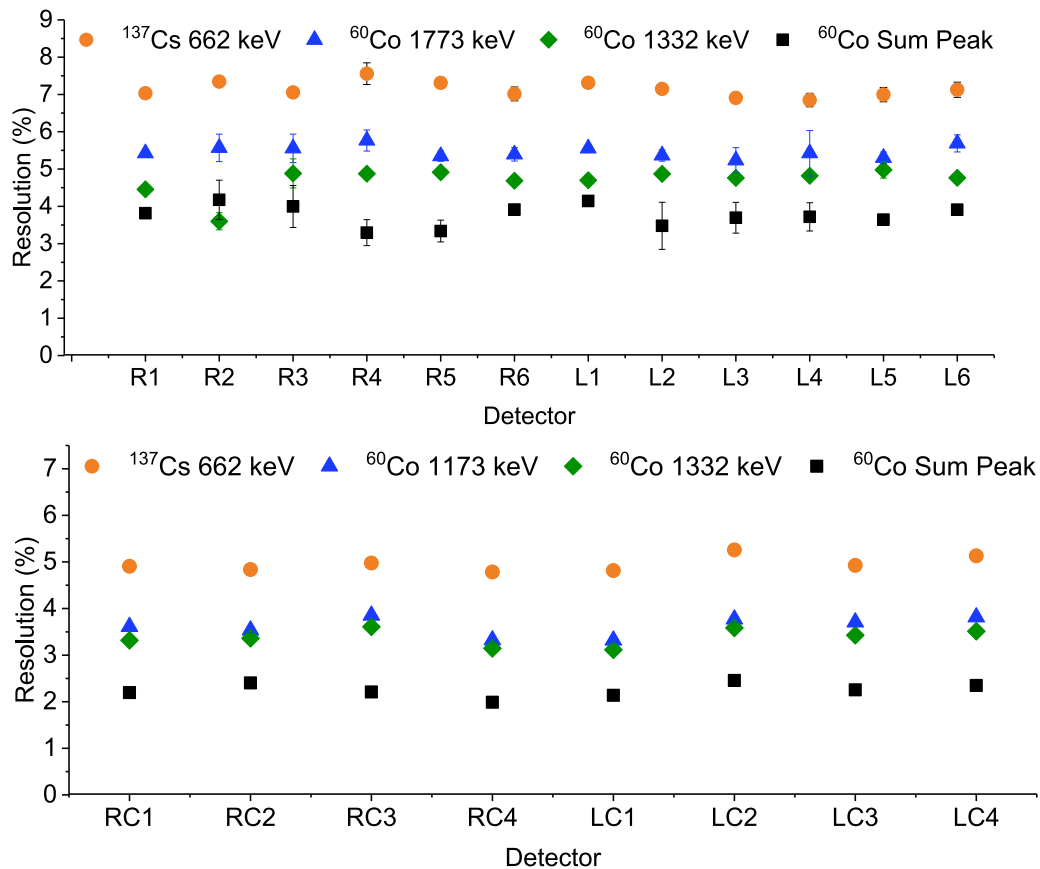


Fig. 4. The energy resolution (%) of the new individual SuN++ segments obtained for the ⁶⁰Co and ¹³⁷Cs sources. The top panel shows the energy resolution for the NaI(Tl) detectors and the bottom panel shows the energy resolution for the CeBr₃ detectors. If error bars are not visible they are within the size of the data points.

for each event, related to the number of γ rays involved in a γ -cascade that de-excites from an excitation energy level in the child nucleus. The spectra registered by a large NaI(Tl), small NaI(Tl), and CeBr₃ segment are shown in black in panels (d), (e), and (f) in Fig. 5. In all panels, the experimental spectra are shown in black and the Geant4 simulated spectra are shown in red, where excellent agreement between the Geant4 simulations and the experimental spectra is observed.

A summing efficiency of 51.1(1)% is obtained for a ⁶⁰Co source located at the center of SuN++. The variation of summing efficiency as a function of source position is illustrated in Fig. 6. In β -decay experiments with SuN++, the implantation point for the ions is at the center position of the total absorption spectrometer, and as such additional Geant4 simulations were performed near the center of the

detector to understand the summing efficiency in this region. The efficiency slowly decreases as the source is moved farther away from the center of the detector on both sides, due to the decrease in solid angle covered by the SuN++.

3.2. In-beam commissioning of SuN++

The commissioning of SuN++ was performed using a ⁷⁰Cu beam produced at FRIB at Michigan State University, behind the Low Energy Beam and Ion Trap (LEBIT). A ⁸²Se primary beam of energy 215 MeV/u was impinged on a 4.88 mm thick ¹²C target. The ⁷⁰Cu ions were selected by the Advanced Rare Isotope Separator (ARIS) [37], and then passed through a momentum compression beam line and into the

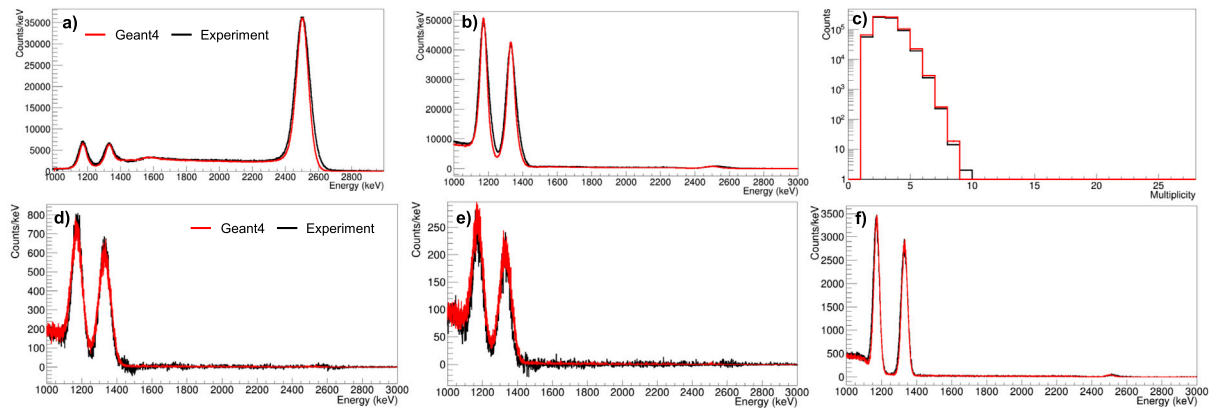


Fig. 5. Various spectra from a ^{60}Co γ -source compared to Geant4 simulations of SuN++. Black lines show the experimental data, and Geant4 simulations in red lines. The top panel includes the (a) total absorption, (b) sum of segments, and (c) multiplicity spectra. The bottom panel shows (d) an individual large NaI(Tl) segment, (e) small NaI(Tl) segment, and (f) CeBr3 segment.

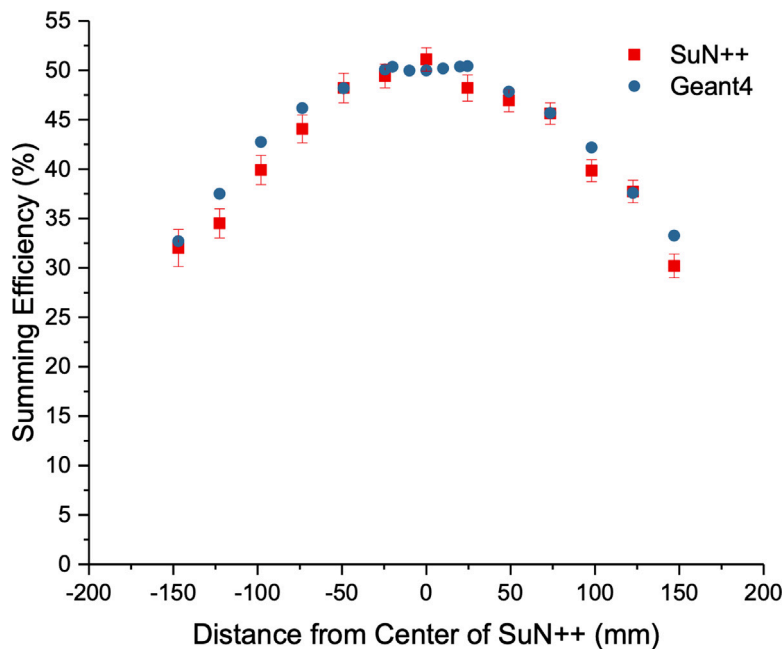


Fig. 6. The summing efficiencies (%) calculated from experimental data and Geant4 simulations as a function of position, where the center of SuN++ is at position 0 mm. The experimental data is shown in the red squares and the Geant4 simulations are shown in the blue dots. The error bars associated with the data points are within the size of the data points if they are not visible.

gas stopping cell where they were thermalized [38]. ^{70}Cu ions were selected in a dipole mass separator with resolution ≈ 1500 and transmitted to the LEBIT 9.4 T Penning trap mass spectrometer at FRIB [39]. The time-of-flight ion-cyclotron-resonance (ToF-ICR) [36,40] technique was used to detect and quantify the percentage of ground ($J^\pi = 6^-$) and two isomeric ($J^\pi = 3^-, 1^+$) states in the ^{70}Cu beam, as shown by the ToF-ICR spectrum in Fig. 7.

In order to isolate the β decay from the $J^\pi = 6^-$ ground-state of ^{70}Cu within LEBIT, the ions of the two isomers were selectively driven to large radial amplitudes via targeted dipole excitation and removed [41]. The efficiency of this procedure is greater than 90%, based on offline tests using pure samples of $^{39}\text{K}^+$. The isolated ground-state ions were transmitted to SuN++, where they were implanted on a double-sided silicon strip detector (DSSD) installed downstream of LEBIT. The DSSD consists of 16 horizontal and 16 vertical strips with a width of 1.2 mm. The DSSD output was fed into high-gain preamplifiers, and the β -decay electrons following ^{70}Cu ion implantation were detected with the DSSD. The detected β -decay electrons in the DSSD

within threshold were used as a β -tag to generate γ -ray spectra with SuN++. The isolation of the ground-state of ^{70}Cu was confirmed with the β -delayed γ -rays observed in SuN++. The child nucleus of ^{70}Cu (^{70}Zn) is stable.

The total absorption, sum of segments, and multiplicity spectra of ^{70}Zn populated via the β decay of the ^{70}Cu ground-state are shown in the left, middle, and right panels of Fig. 8, respectively. Known levels in ^{70}Zn previously reported by Van Roosbroeck et al. were simulated in Geant4 to account for the detector response of SuN++ [43]. The levels in ^{70}Zn populated via the ground-state β decay of ^{70}Cu are known up to 5.06 MeV with the β -decay Q-value (Q_β) being 6.584(3) MeV. To account for the apparent feeding shown in Fig. 8 above the last known level in ^{70}Zn , pseudo-levels were created using RAINIER [44]. RAINIER relies on select models for the nuclear level density, γ -ray strength function, and spin distribution, as well as known level schemes at low energies. RAINIER is used to simulate decays from a given excitation energy level in the nucleus of interest. The levels simulated with RAINIER are pseudo-levels, which are representative of an energy

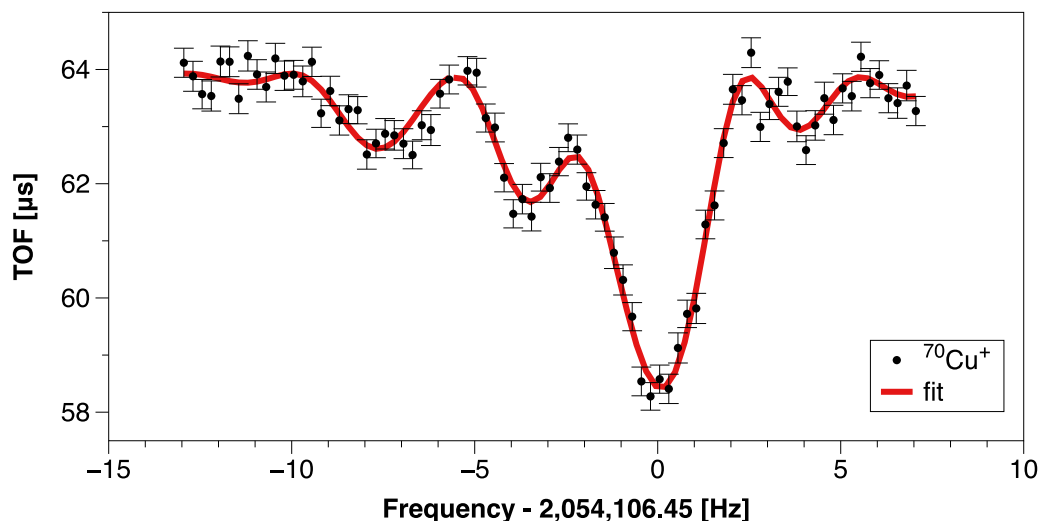


Fig. 7. Time-of-Flight Ion Cyclotron Resonance (ToF-ICR) [36] scans from LEBIT when ^{70}Cu was delivered to the Penning trap during the SuN++ commissioning experiment. The figure shows the mix of ground and isomeric states in ^{70}Cu .

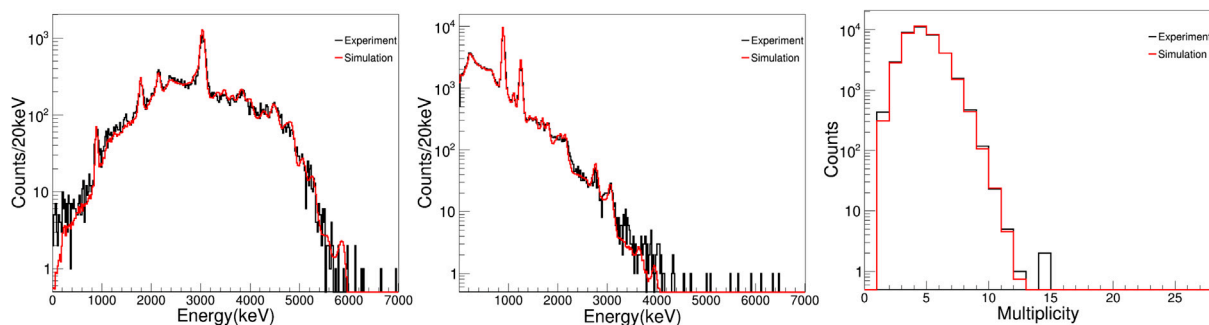


Fig. 8. The total absorption spectrum, sum-of-segments, and multiplicity in ^{70}Zn resulting from the ground-state decays of ^{70}Cu .

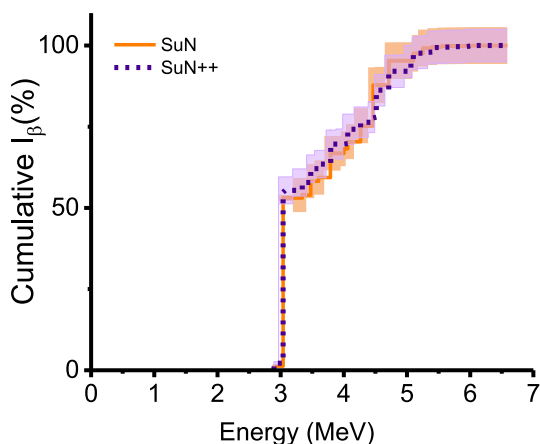


Fig. 9. The cumulative β -feeding (I_β) from the β decay of the $J^\pi=6^-$ ground state of ^{70}Cu into ^{70}Zn . The pink band shows the cumulative I_β from SuN++, compared with the cumulative I_β from SuN in orange [42].

band that corresponds to the resolution of SuN++. Pseudo-levels in RAINIER were simulated from the highest-known level to the Q_β and then fed through the SuN++ Geant4 simulations to account for detector response. A χ^2 minimization procedure was used to simultaneously fit simulated data to all three experimental spectra, and is shown by the red line in Fig. 8; this procedure is similar to those outlined by Dombos

et al. and Lyons et al. [24,25]. Fig. 8 shows good agreement between the fits and the experimental data.

From the best fits of χ^2 -minimization, the I_β distribution was extracted. The cumulative I_β (%) as a function of excitation energy is shown in Fig. 9. The I_β values were compared to the I_β obtained from the β -decay of the ground state of ^{70}Cu with SuN [42]. The measurement by Ronning et al. [42] (orange error band) agrees with the current measurement (pink error band) within error. Although the cumulative I_β values extracted with each detector agree within error, the shape of the β -feeding curves varies slightly between SuN and SuN++. This is likely the result of the improved resolution of SuN++, allowing for better distinction between peaks in the total absorption spectrum and more distributed β -feeding to each excitation energy.

Understanding the multiplicity response of SuN++ is important, as both the original SuN segments and the new SuN++ segments cover restricted azimuthal angles. The multiplicity response of the SuN++ detector to multiplicity (M) = 1, 2, and 3 events is shown in Fig. 10 a. The experimental multiplicities are shown by the colored solid lines, with the Geant4 multiplicity shown by the dashed lines. The $M = 1$ response (magenta) comes from the ^{137}Cs γ -source, the $M = 2$ response (orange) comes from the ^{60}Co γ -source and the $M = 3$ response (blue) comes from the ^{70}Cu β -decay data. The state in ^{70}Zn with the highest I_β feeding percentage from ^{70}Cu is at 3038.30(16) keV, which decays by emitting three γ -rays: 1251.7, 901.7, and 884.88 keV. Thus, the peak in the total absorption spectrum at 3038.30(16) keV corresponds to $M=3$ events. The experimental $M = 3$ spectrum was produced by gating on this TAS peak to produce a multiplicity spectrum, and compared to Geant4 simulations of this level. The Geant4 simulations are in excellent agreement with experimental data.

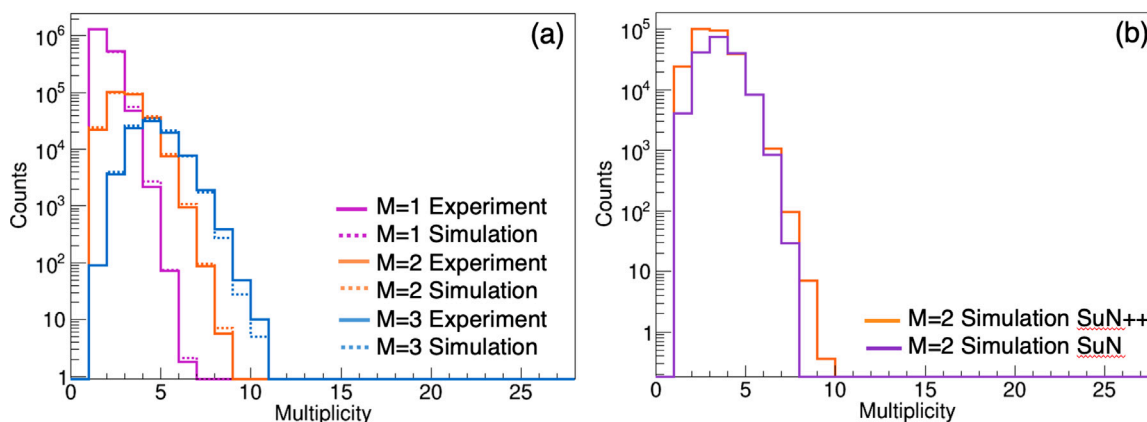


Fig. 10. Panel (a) shows the segment multiplicity spectra from SuN++ for γ -ray multiplicity 1 (magenta), 2 (orange), and 3 (blue) events in solid lines compared to Geant4 simulations (dashed lines). Panel (b) shows the segment multiplicity spectra from Geant4 simulations of a ^{60}Co source ($M = 2$) in SuN++ (orange) vs SuN (purple).

Fig. 10 b shows the multiplicity response of SuN (purple line) and SuN++ (orange line) to a ^{60}Co source simulated in Geant4. The highest multiplicity that can be registered in SuN is $M = 8$ and the highest multiplicity that can be registered in SuN++ is $M = 28$. We can see in Fig. 10 b that although the multiplicity responses of the two total absorption spectrometers are not identical, they are very similar, despite the large variation in azimuthal angles of the SuN++ detectors.

4. Conclusions

In this paper, we present the details of the upgrade of SuN to the SuN++ total absorption spectrometer. The summing efficiency of the SuN++ detector was tested using ^{60}Co and ^{137}Cs sources and was found to be 51.1(1)% at 2.5 MeV at the center of SuN++, compared to 65(2)% from SuN [23]. The results obtained from measurements for all detector types were found to be in excellent agreement with GEANT4 simulations. SuN++ was commissioned at FRIB behind LEBIT which delivered isomer-separated beams of ^{70}Cu . The experimentally obtained total absorption, sum-of-segments, and multiplicity spectra were also in good agreement with the simulation for β decay of ^{70}Cu to ^{70}Zn , and with previous results by Ronning et al. [42]. SuN++ now allows for better energy resolution and identification of populated states through β decay using the newly added CeBr_3 segments, which have an average resolution at 1.332 MeV of 3.38(4)%, compared to 5.8(2)% from the SuN segments. The additional 20 segments in the SuN++ total absorption spectrometer allow for the detection of higher-multiplicity cascades, confirmed by the multiplicity response illustrated in Fig. 10. These upgrades, along with the planned addition of electron detectors to SuN++, are ideal for the investigation of indirect measurements of (n, γ) reactions relevant to s -, i - and r -process nucleosynthesis.

CRediT authorship contribution statement

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Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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Data availability

Data will be made available on request.

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